



### Abstract

Asymmetries in plasma density irregularity generation between the leading and trailing edges of large-scale plasma density structures in the high-latitude ionosphere are investigated. A model is developed that evaluates the gradient-drift instability (GDI) growth rate differences across the gradient reversal that is applicable at all propagation directions and for a broad range of altitudes spanning the entire lower ionosphere. In particular, the model describes asymmetries that would be observed by an oblique scanning radar near density structures in the polar cap such as polar patches. The dependencies on the relative orientations between the directions of the gradient reversal, plasma convection, and wave propagation are examined at different altitudinal regions. At all altitudes, the largest asymmetries are expected for observations along the gradient reversals, e.g. when an elongated structure is oriented along the radar boresite. The convection direction that results in the strongest asymmetries exhibits a strong dependence on the altitude, with the optimal convection being parallel to the gradient reversal in the E region, perpendicular to it in the F region, and at some angle between these extremes in the transitional region.

### Introduction

The gradient-drift instability (GDI) allows for wave growth when a drifting plasma has a density gradient. Observations have found asymmetries between the leading and trailing edges of large-scale density structures in the ionosphere (e.g. Weber et al., 1984; Milan et al., 2002). GDI linear theory can be applied to the field of view (FoV) of a HF coherent scatter radar, such as in the Super Dual Auroral Radar Network (SuperDARN), to provide some insight into this phenomenon.



Height dependence of the anisotropy parameter,  $\psi$ , and the conductivity ratio, R. Both decrease by several orders of magnitude between E-region and F-region altitudes.





- Height of SuperDARN HF radar beam.
- Black-white line represents large-scale density gradient reversal
- Uniform field of gradient magnitudes
- Gradient direction reverses at the line
- Gradient reversal direction:  $\theta$
- Convection direction:  $\phi$
- Vary both  $\theta$  and  $\phi$  independently Large  $\gamma \rightarrow$  greater structuring  $\rightarrow$  more
- observed backscatter

SuperDARN FoV for a particular gradient reversal and convection direction.

- Different orientations produce different degrees of contrast in the growth rate
- Asymmetry in growth rate implies asymmetry in observed backscatter
- Need to quantify contrast over  $\bullet$ entire FoV

# **Asymmetry in Plasma Irregularity Growth** Near Large-Scale Gradient Reversals

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Slant Range (km) Different slant ranges sample the E region, F region, and a transitional region between the two.

- GDI growth rate at each cell within the McMurdo Station

# **Methods for Quantifying Asymmetry**

# Numerical

- 1.Consider a series of points within the FoV that are mirrored over the gradient reversal
- 2. Find wavevector  $(\vec{k})$  and altitude (*R*) for virtual beam
- 3. Find the growth rate at each point
- 4.Subtract trailing edge point  $\gamma$ from corresponding leading edge point  $\gamma$  for  $\Delta \gamma$
- 5. Average over all point pairs

## Results





- R decreases by several orders of magnitude
- $\Delta\gamma$  smoothly transitions from E region to F region pattern
- Low sensitivity to changes in altitude above 130 km



# Analytic

- 1.Consider a point on the gradient reversal within FoV
- 2. Wavevector,  $\vec{k}$ , defined by angle with boresight,  $\alpha$
- 3. Find  $\Delta \gamma$  in terms of  $\alpha$ ,  $\theta$ , and  $\phi$
- 4.Integrate over FoV to find expression for average  $\Delta \gamma$

 $\Delta \gamma = H(\theta, \varphi) [\sin(\varphi - \theta) - C \sin(\varphi + \theta) + R(\cos(\varphi - \theta) - C \cos(\varphi + \theta))]$  $H(\theta, \varphi) = \Theta(\varphi - \theta + \pi) - \Theta(\theta - \varphi - \pi) + 2\Theta(\varphi - \theta - \pi) - 2\Theta(\varphi - \theta)$ 

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