

Parametric Wave Growth in a Hybrid PIC/Fluid Simulation of the Equatorial E Region.

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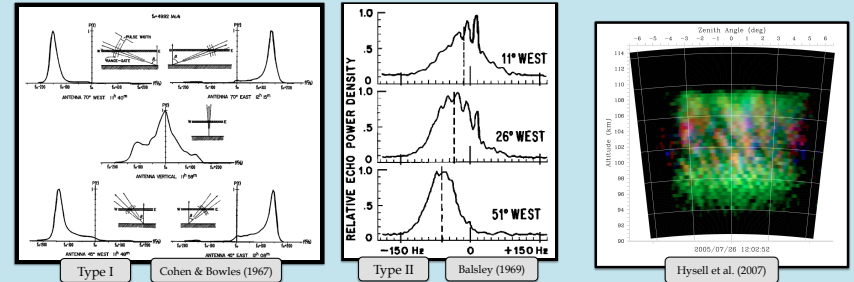


Background

This work models the **equatorial ionosphere** around 100 km. The plasma is **electrostatic** and **quasineutral**. Ions are demagnetized by collisions while electrons $E \times B$ drift westward. Density gradients and electric fields can cause this plasma to become unstable to both the **gradient-drift** and **Farley-Buneman** instabilities. These are the **first simulations** of coupled Farley-Buneman/gradient-drift instability initiated by a large-scale density perturbation.

Motivation

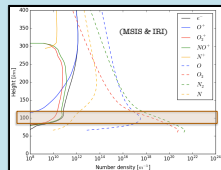
The 50-MHz radar at Jicamarca has observed so-called Type I and Type II irregularities for decades (e.g. Cohen & Bowles (1967) and Balsley (1969)). Type I spectra are narrow with a constant Doppler shift; Type II spectra are broad with Doppler shift that increases with look angle. Type I are attributed to the Farley-Buneman (FB) instability and Type II are attributed to the gradient-drift (GD) instability, but both are special cases of the same dispersion relation. More recent observations by Hysell et al. (2007) suggest that meter-scale irregularities observed at Jicamarca can trace kilometer-scale structures. Thus far, no simulations have reproduced this coupled turbulence.



Simulation

- Ions: particle-in-cell method with ≈ 1 billion particles.
- Electrons: inertialess, isothermal fluid (quasineutral with ions).
- Density: initialized as a Gaussian (FWHM = 146 m) perturbed by a 512-m wave.
- Electrostatic potential: Calculated from background electric field ($E \times B$ drift), total density, ion flux, and electron fluid quantities.
- Background densities and temperatures taken from MSIS & IRI model data (01 Jan 2000 @ noon).

Physical Setup



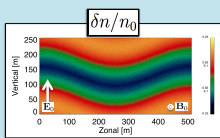
Equatorial E region around 100 km
Dominant neutrals are N_2 and dominant ions are NO^+ . Ions are demagnetized by collisions with neutrals.

$(dx, dy) = (0.25, 0.25)$ m
 $(nx, ny) = (2048, 1024)$ cells
 $dt = 5 \times 10^{-5}$ s
 $nt = 8192$ steps
 $B_0 = -2.5 \times 10^{-5} \hat{y}$ T
 $m_e = 9.1 \times 10^{-31}$ kg
 $m_i = 5.0 \times 10^{-26}$ kg
 $m_n = 4.6 \times 10^{-26}$ kg
 $n_0 = 10^{10} m^{-3}$
 $\nu_e = 3.0 \times 10^4 s^{-1}$
 $\nu_i = 3.0 \times 10^3 s^{-1}$
 $T_e = T_i = T_n = 220$ K

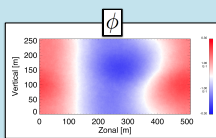
Potential Solver

$$\epsilon \equiv \begin{pmatrix} 1 & -\kappa_e \\ \kappa_e & 1 \end{pmatrix} \quad \kappa_e \equiv \frac{\Omega_e}{\nu_e} \leftarrow \text{Cyclotron frequency} \quad \nu_e \leftarrow \text{Collision frequency}$$

$$\nabla \cdot (n \epsilon \nabla \phi) = \nabla \cdot \left\{ \epsilon \left[n E_0 + \left(\frac{\gamma T_e}{e} \right) \nabla n \right] + (1 + \kappa_e^2) \frac{m_e \nu_e}{e} n v_i \right\}$$

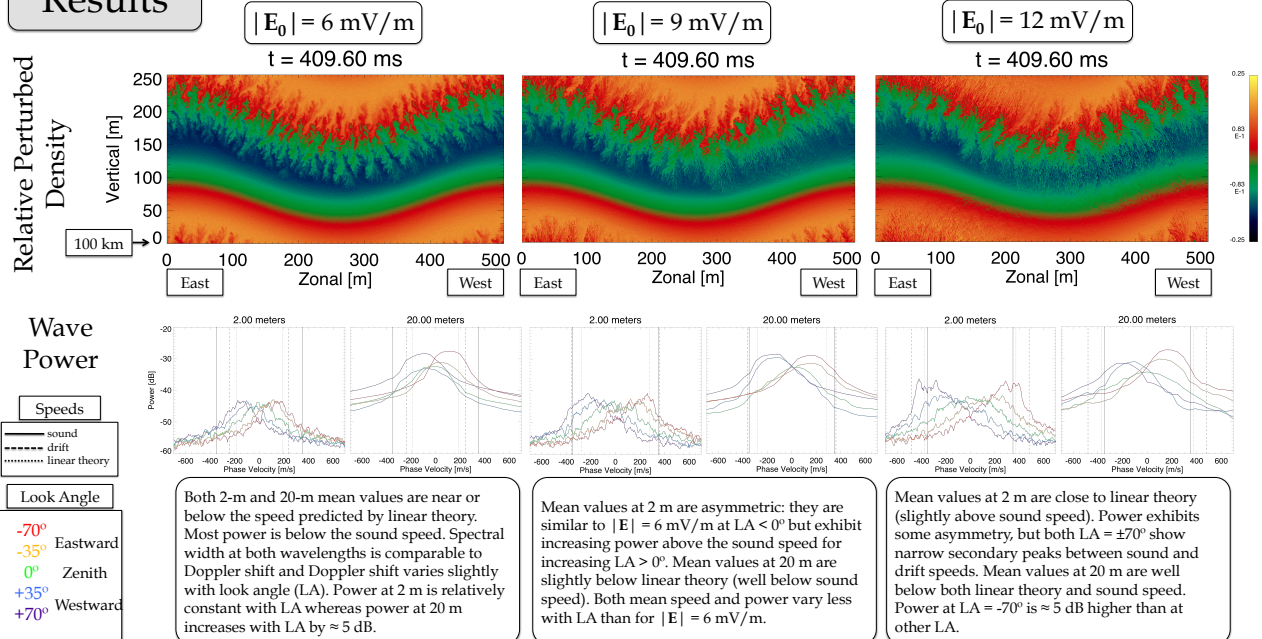


Initial relative perturbed density



Initial electrostatic potential

Results



Conclusions

Two-meter irregularities with speeds at or above the sound speed grow asymmetrically with increasing $|E|$.

Power spectra are mostly Type II with subsonic mean but show fast, narrow peaks at 2 m for largest $|E|$.

Fourier spectra (not shown) for $|E| > 9 \text{ mV/m}$ exhibit an initial burst of meter-scale power before growth of longer wavelengths