



*The Johns Hopkins University
Applied Physics Laboratory
Geospace Remote Sensing*

Coherent Radar Studies of the Ionosphere

by

**Raymond A. Greenwald
The Johns Hopkins University
Applied Physics Laboratory**

June 21, 1992



Contents

- Characteristics of Coherent Scatter Cross-Sections
- Siting and Propagation Considerations
- Examples of Coherent Scatter Radars
- Plasma Instability Processes
- Spectral Characteristics of E- and F-Region Backscatter
- The Drift Motion of Ionospheric Irregularities
- Applications of Coherent Radars to Solar-Terrestrial Research
 - Convection Dynamics
 - Gravity Wave Excitation and Propagation
- SuperDARN



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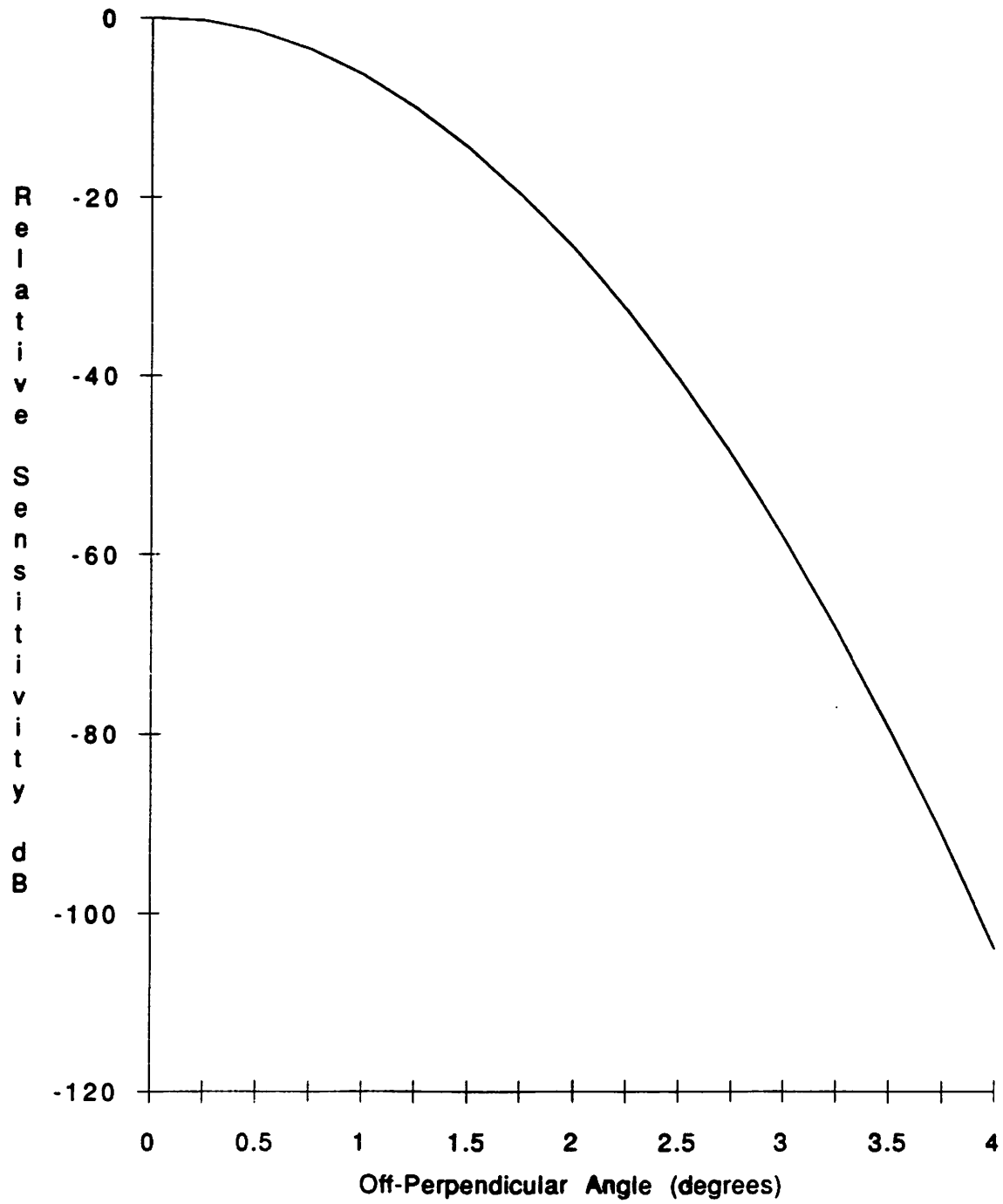
Incoherent Scatter Cross-Section

$$\sigma_B = 4 \pi r_e^2 N / (1 + T_e / T_i)$$

Where:

N:	Electron density
T_e:	Electron Temperature
T_i:	Ion Temperature
r_e:	Classical radius of electron

SENSITIVITY DEPENDENCE FOR ASPECT RATIO (70) IRREGULARITIES



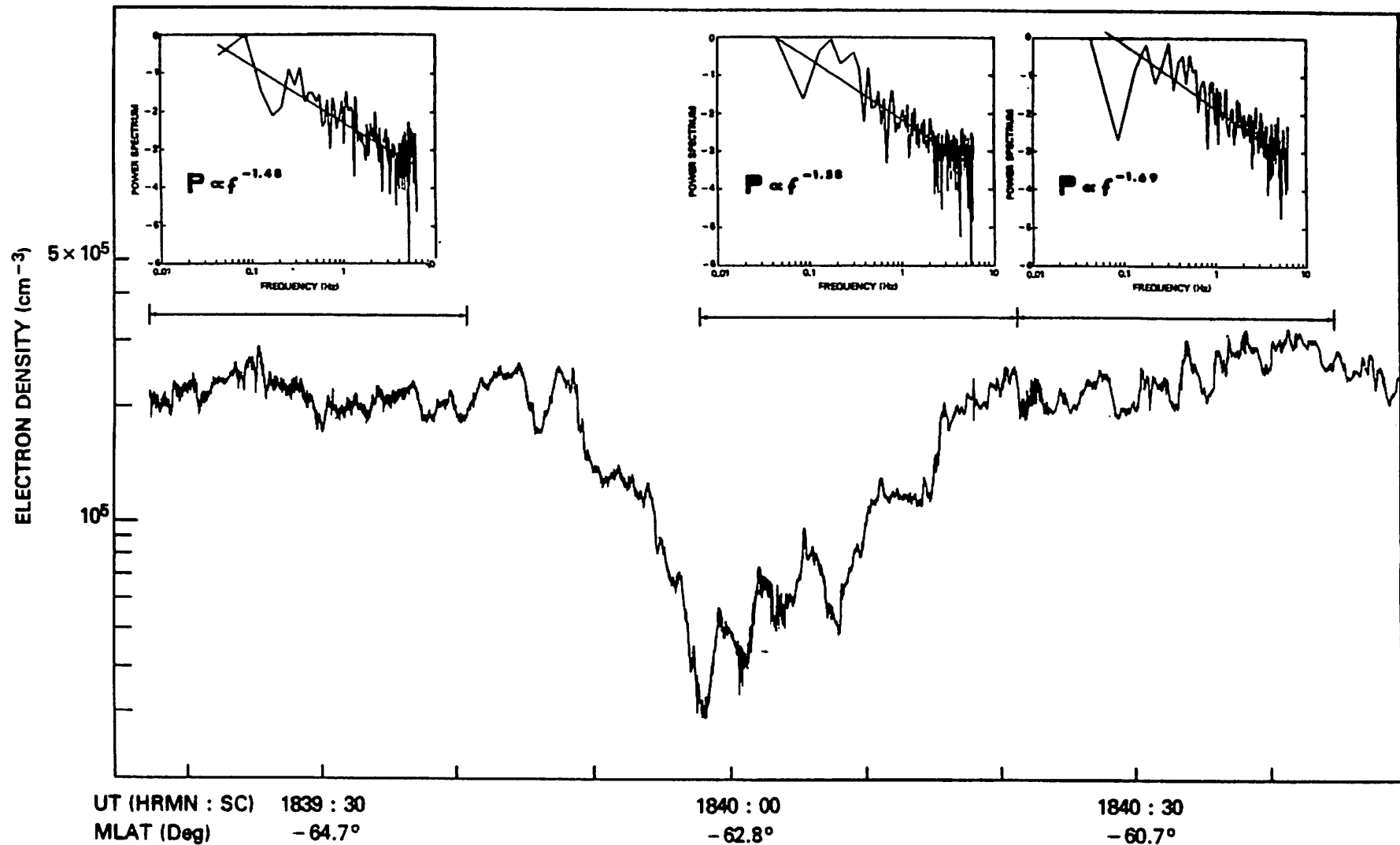


FIGURE 2. An expanded time-scale plot of electron density irregularities in the diffuse aurora of Figure 1. The power spectra indicate the frequency dependence of the least-squares power law. Frequency f is directly proportional to wave-number k , assuming the irregularities are spatial structures.

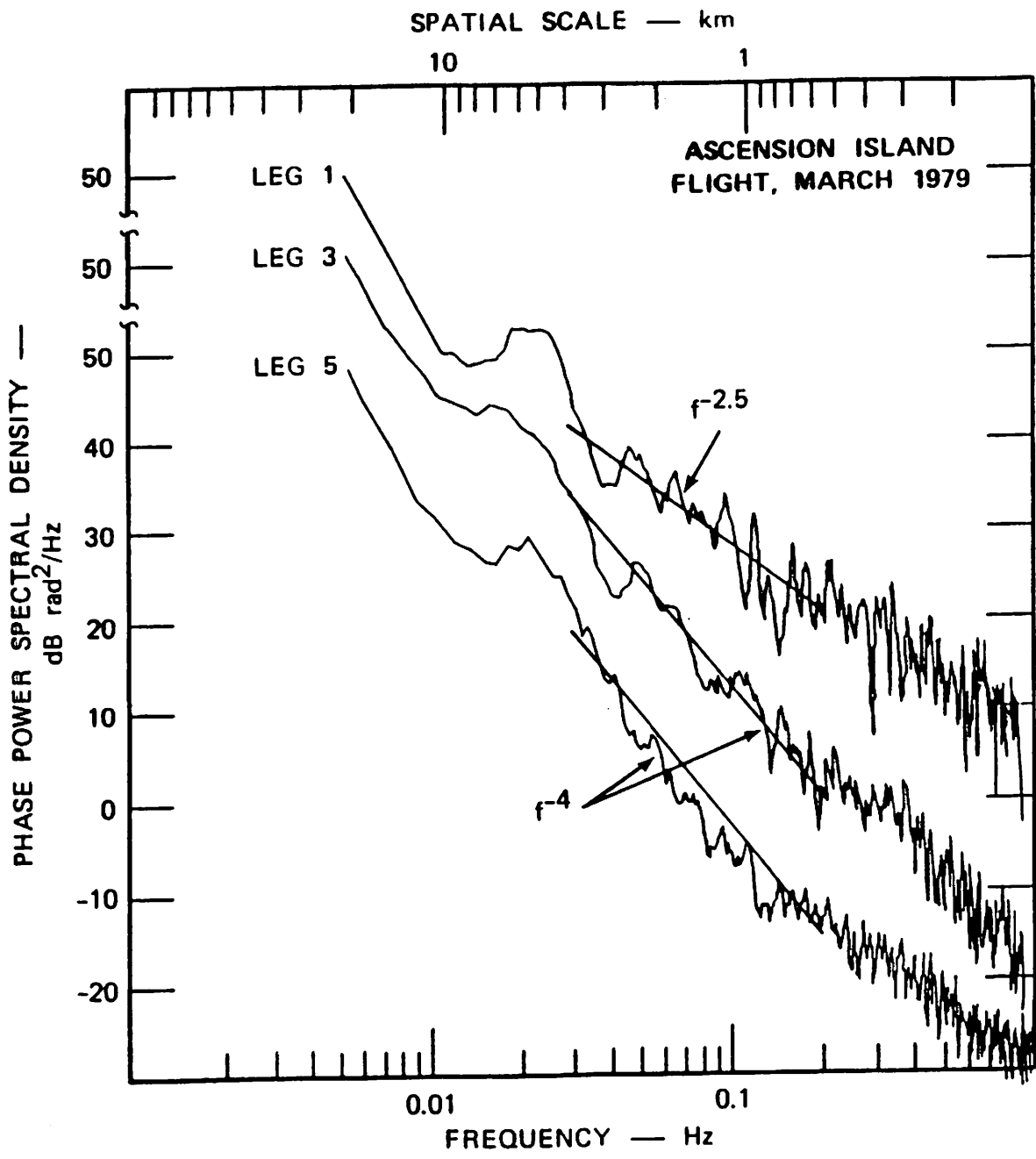


FIGURE IV.A.5 PHASE POWER SPECTRA FOR THE BUBBLE



Comparison of Coherent Scatter and Incoherent Scatter Cross-Sections

For $n = 3$,

$$\mathbf{P}_3(2k, 0, 0) \Rightarrow \Pi^{3/2} \alpha / 4 \lambda_0 k^4$$

and

$$\sigma_c / \sigma_{inc} = \langle \Delta N^2 \rangle (1 + T_e / T_i) \Pi^{1/2} \alpha / (16 N \lambda_0 k^4)$$

Assume

$$\lambda_0 = 10 \text{ km}$$

$$k = \Pi / 10 \quad (15 \text{ MHz})$$

$$a = 70$$

$$N = 5 \times 10^5$$

$$\langle \Delta N^2 \rangle^{1/2} / N = 10\%$$

$$T_e = T_i$$

Then: $\sigma_c / \sigma_{inc} \approx 160$

MEASUREMENT OF FLUCTUATION LEVEL

BACKSCATTER CROSS SECTION (BOOKER, 1950)

$$\sigma_B = r_e^2 \langle \Delta N^2 \rangle P_3(k_1, k_2, k_3)$$

WHERE $r_e = 2.8 \times 10^{-15} \text{ m} = \text{CLASSICAL ELECTRON RADIUS}$

$\langle \Delta N^2 \rangle = \text{MEAN SQUARE ELECTRON DENSITY FLUCTUATION LEVEL}$

$P_3(k_1, k_2, k_3) = \text{3D SPATIAL POWER SPECTRUM}$

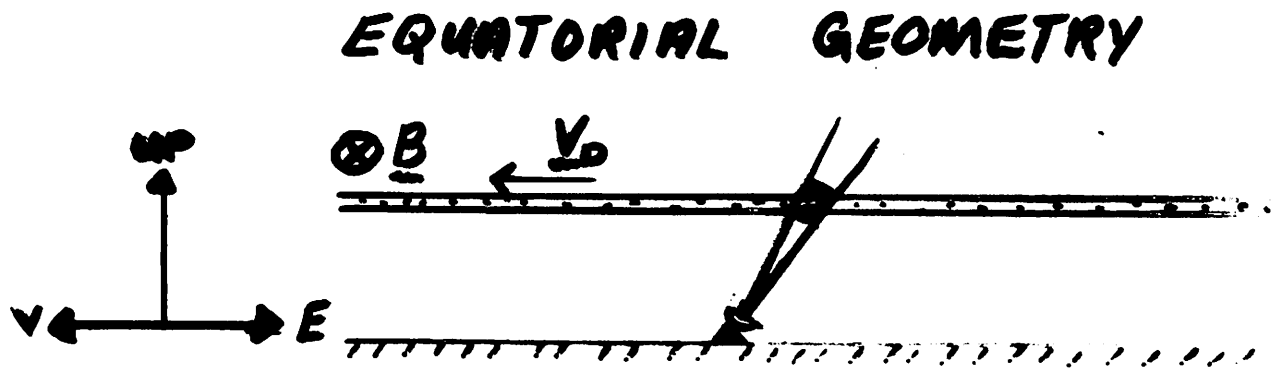
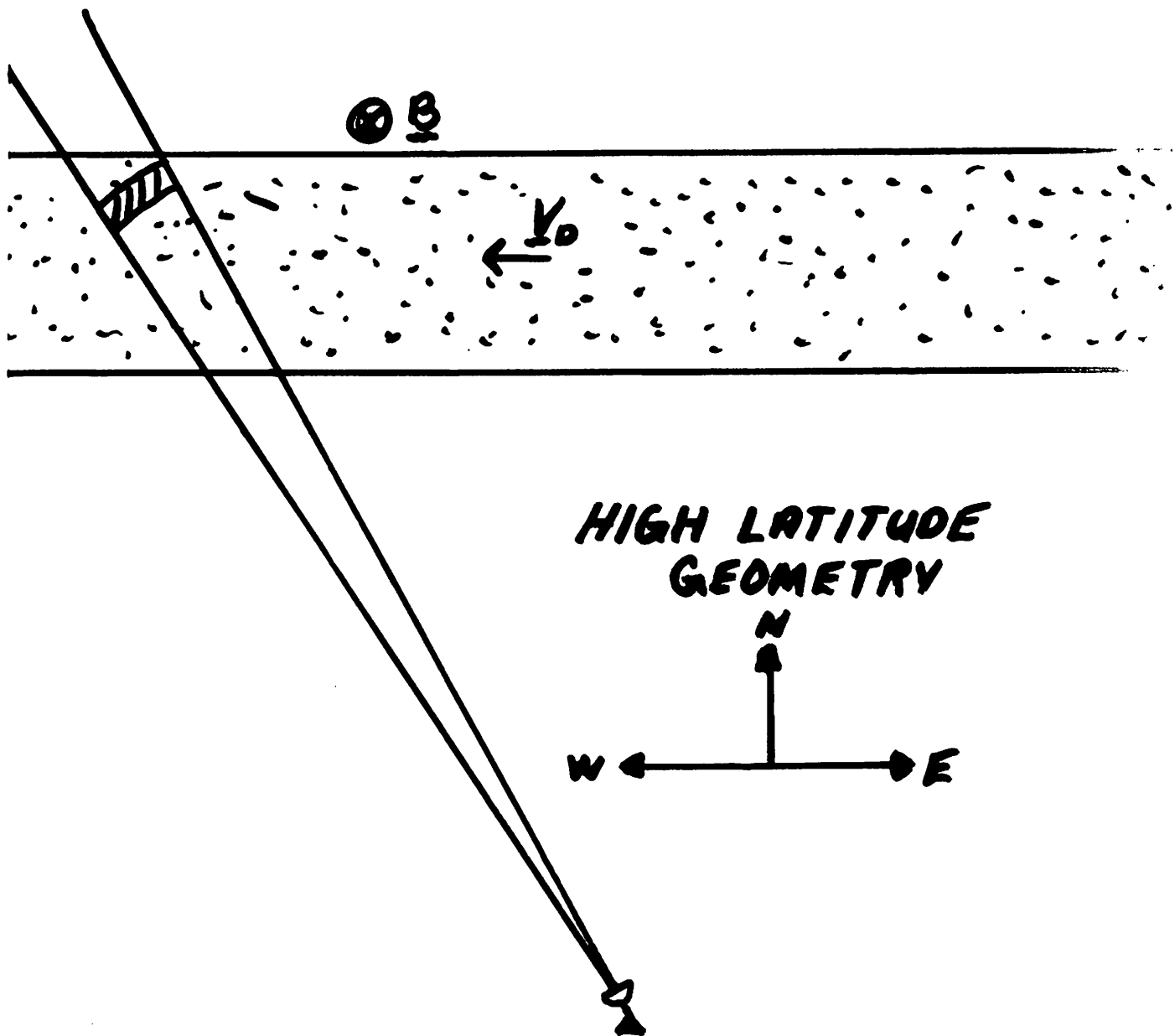
ASSUME

$$P_3(k_{\perp}, 0, k_{\parallel}) = \frac{A_{\perp} e^{-(\lambda_0 \alpha k_{\parallel})^2 (1 + \lambda_0^2 k_{\perp}^2)}}{(1 + \lambda_0^2 k_{\perp}^2)^{n+1/2}}$$

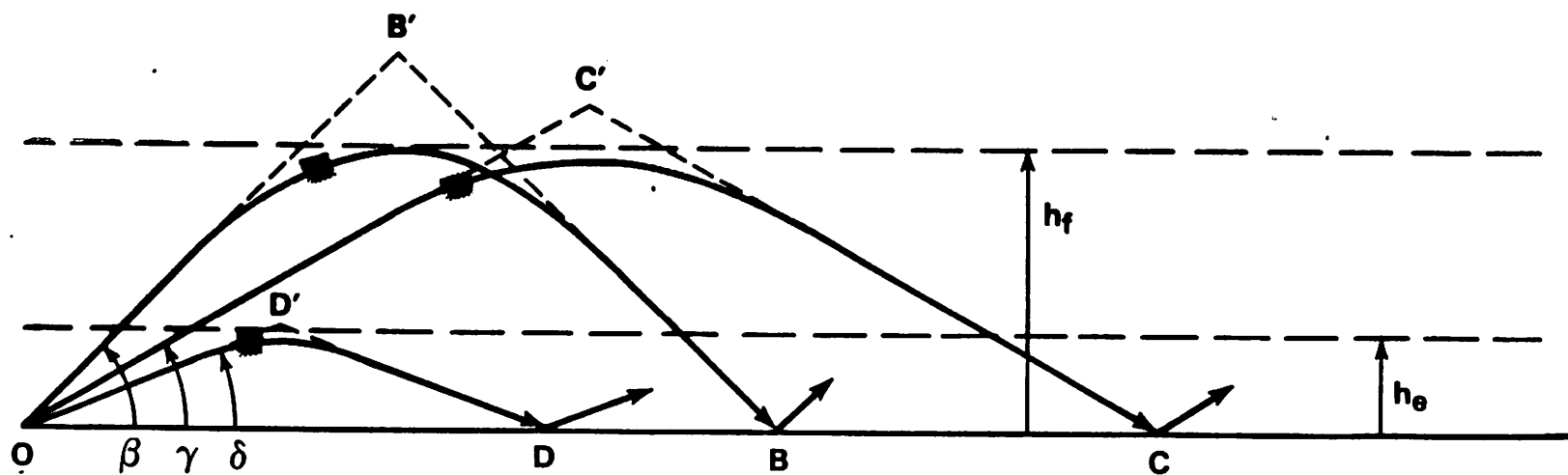
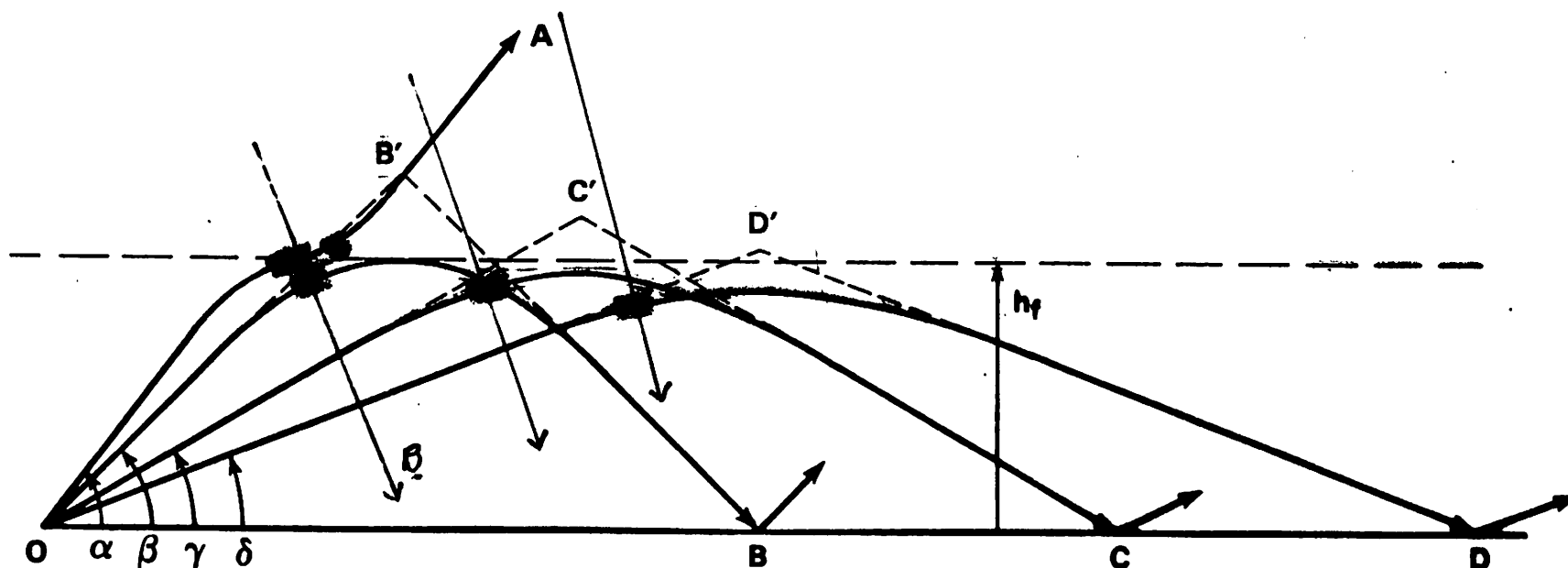
WHERE $\lambda_0 = \text{OUTER SCALE LENGTH OF TURBULENCE}$

$\alpha = \text{ASPECT RATIO OF IRREGULARITIES}$

$$n > 1$$



HF PROPAGATION IN A HORIZONTALLY-STRATIFIED IONOSPHERE





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Drift of F-Region Irregularities

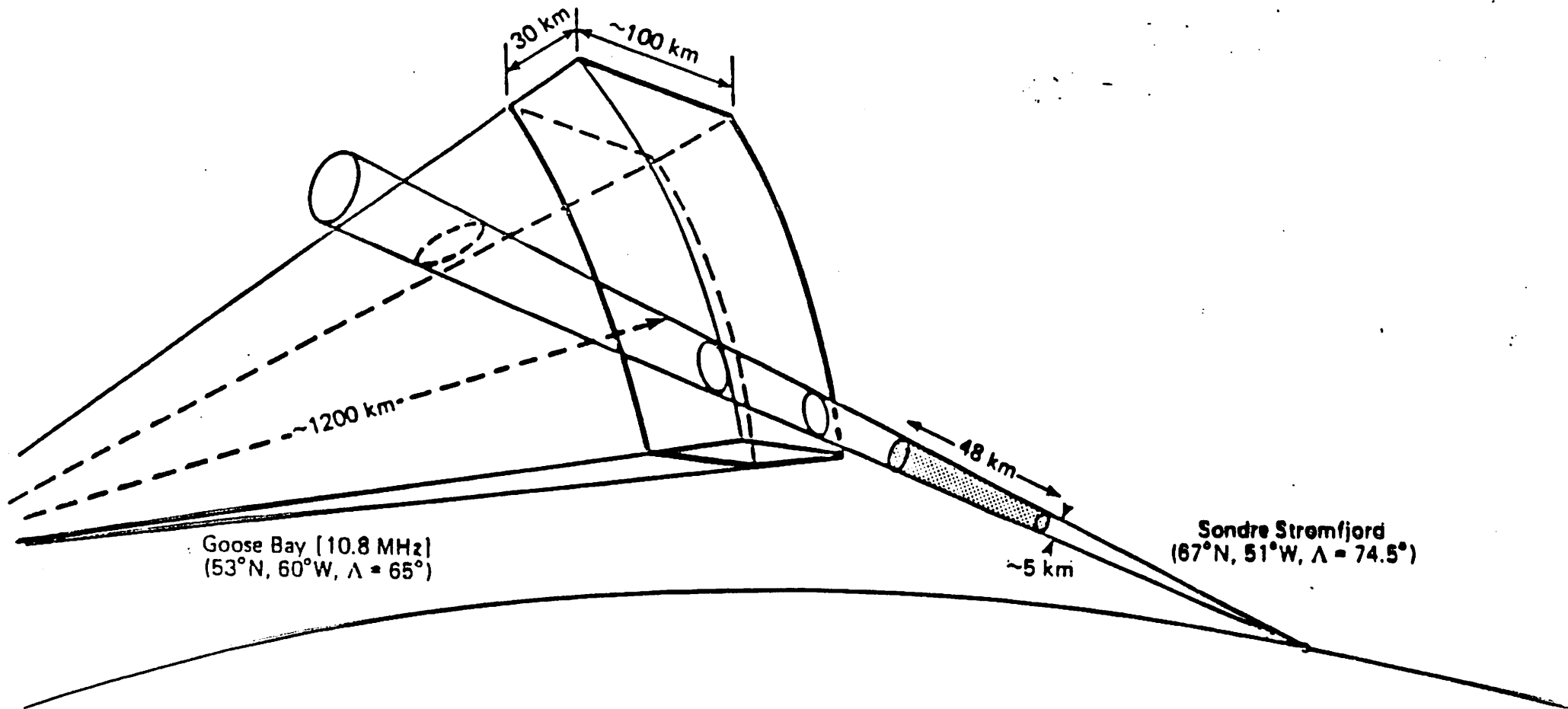
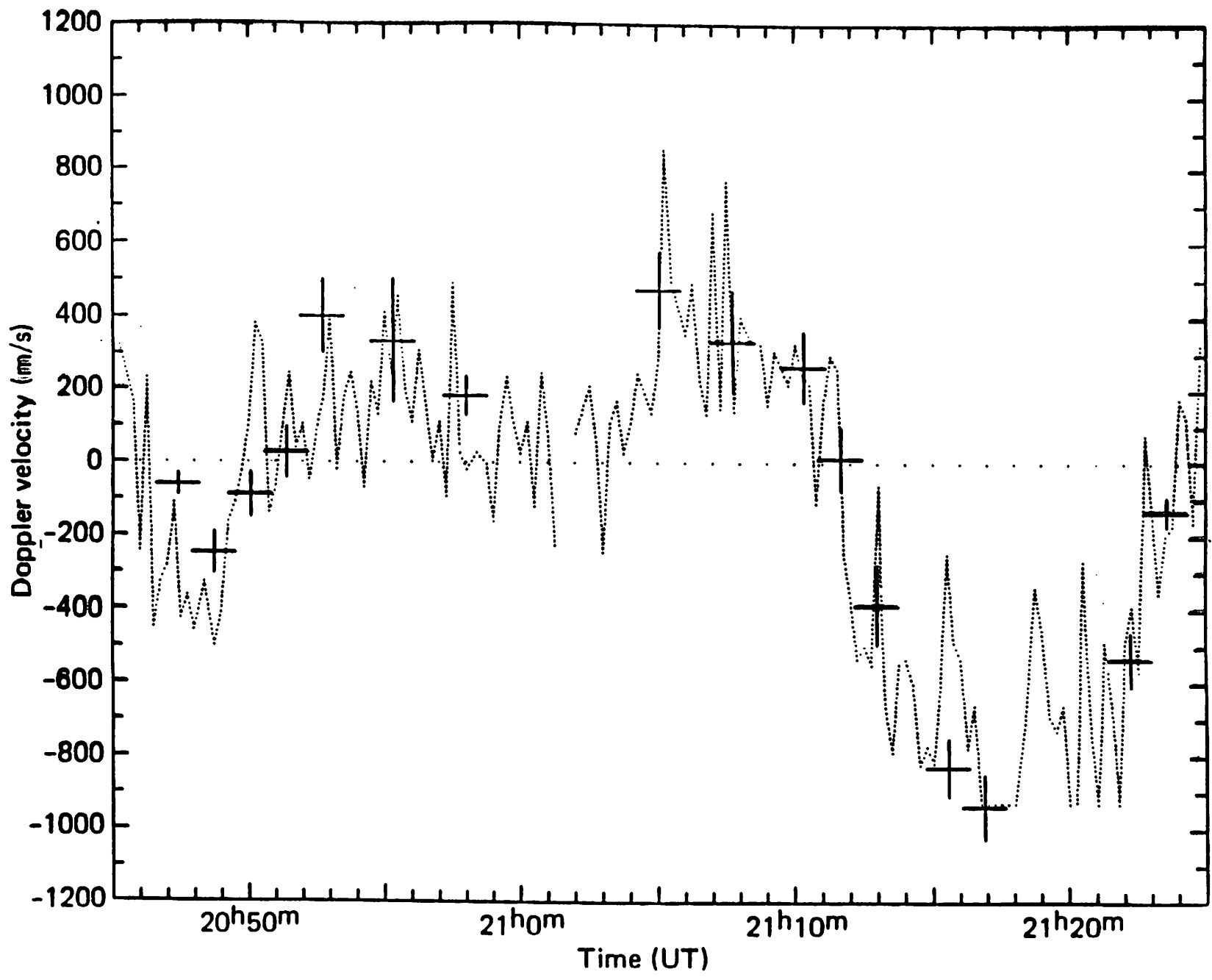


Fig. 1





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Plasma Instabilities



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Categories of Plasma Instability

Source of Free Energy

Streaming of one particle species relative to another

Gradient in plasma density or temperature plus external force

Gradient in plasma density or temperature without external force

Non-Maxwellian particle distribution

Type of Instability

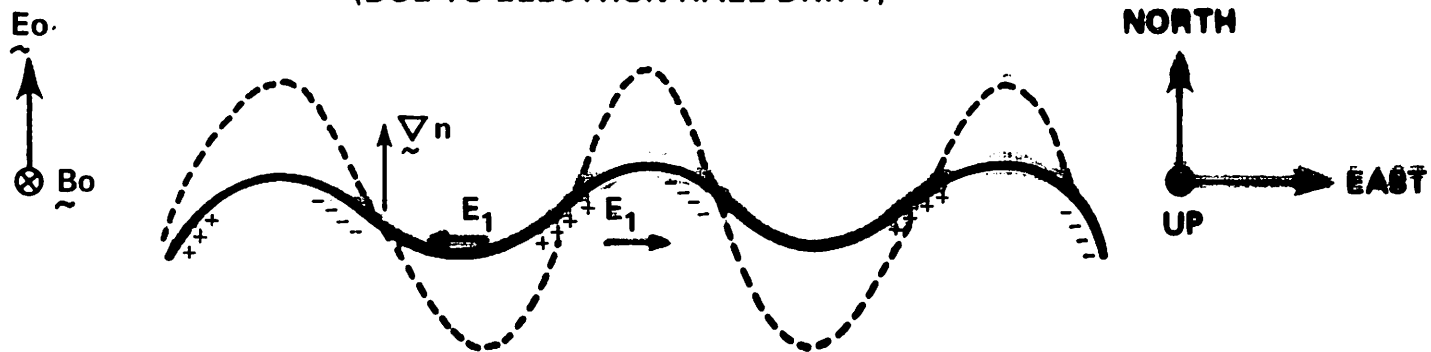
Streaming Instability

Gradient Instability

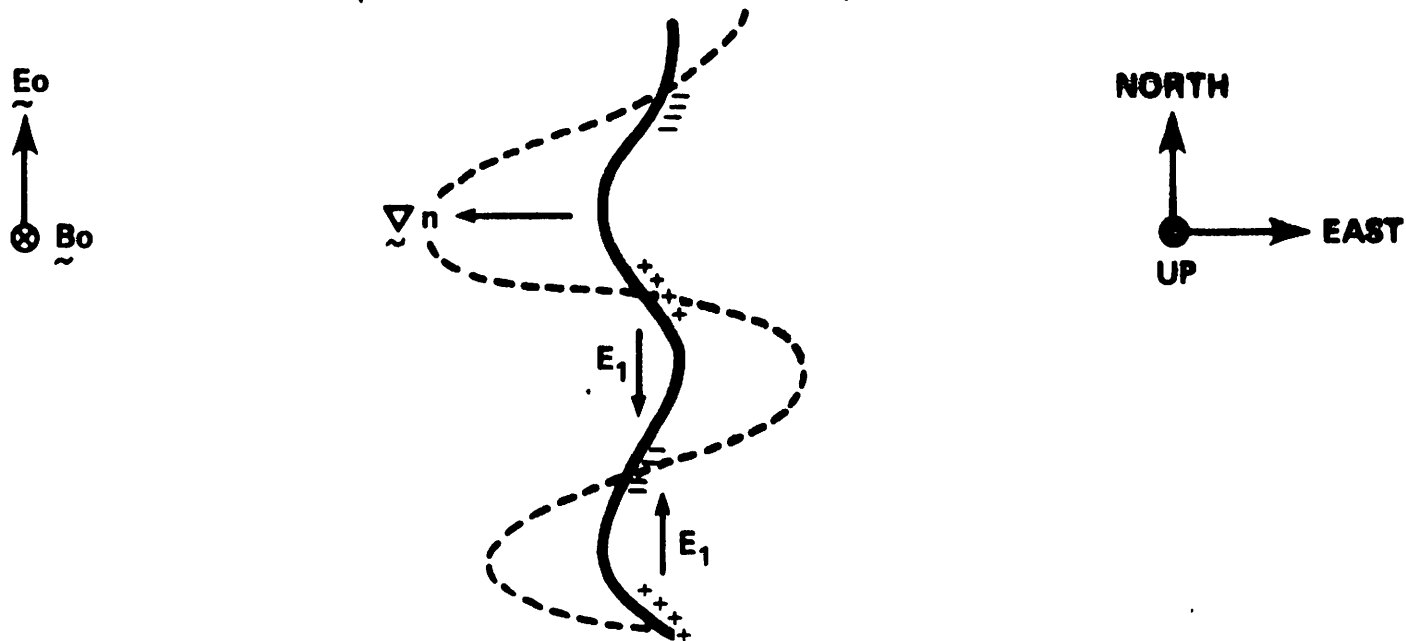
Universal Instability

Kinetic Instability

**E-REGION GRADIENT-DRIFT INSTABILITY
(DUE TO ELECTRON HALL DRIFT)**



**F-REGION GRADIENT-DRIFT INSTABILITY
(DUE TO ION PEDERSEN DRIFT)**



CASE 2. GRADIENT DRIFT INSTABILITY

$$k_{\parallel} = 0 \quad v_d \ll c_s \quad L_N \neq \infty \quad \psi \ll 1$$

$$\Gamma = \frac{k \cdot v_d v_i}{k L_N \Omega_i} - \frac{\psi}{v_i} k^2 c_s^2$$

$$\omega_r = \underline{k} \cdot \underline{v}_d$$

CASE 3. MODIFIED TWO-STREAM INSTABILITY

$$k_{\parallel} \neq 0, \quad L_N = \infty, \quad \psi > 1, \quad \underline{v}_d = \{v_{d\parallel}, v_{d\perp}\}$$

$$\frac{v_i}{\Omega_i} \gtrsim 1$$

$$\Gamma \simeq \frac{1}{v_i} [(k \cdot \underline{v}_d / \psi)^2 - k^2 c_s^2]$$

$$\omega_r \simeq \underline{k} \cdot \underline{v}_d / \psi$$

$$\text{THRESHOLD: } \omega_r / k = c_s$$

CASE 4. ELECTROSTATIC ION CYCLOTRON INSTABILITY

$$k_{\parallel} \neq 0, \quad L_N = \infty, \quad \psi > 1, \quad \underline{v}_d = \{v_{d\parallel}, v_{d\perp}\}$$

$$\frac{v_i}{\Omega_i} \ll 1$$

$$\text{THRESHOLD: } \omega_r / \Omega_i = \left(\frac{\Omega_i^2}{k^2} + c_s^2 \right)^{\frac{1}{2}}$$

E AND F₁ REGION PLASMA STREAMING INSTABILITIES

C_s = ACOUSTIC VELOCITY

V_d = RELATIVE ELECTRON-ION DRAFT $\{V_{d\perp}, V_{d\parallel}\}$

L_N = DENSITY GRADIENT SCALE LENGTH

$$\psi = \left(\frac{N_i V_i}{N_e \Omega_i} \right) \left(1 + \frac{R_e^2 k_y^2}{N_e^2 k^2} \right)$$

$$k_y = 0 \Rightarrow \psi = \left(\frac{N_i V_i}{N_e \Omega_i} \right) \ll 1$$

$$k_y \neq 0 \Rightarrow \psi = \left(\frac{V_i R_e}{N_e \Omega_i} \right) \frac{k_y^2}{k^2} > 1$$

CASE 1. CROSS FIELD TWO-STREAM INSTABILITY
 $k_y = 0$ $L_N = \infty$ $\psi \ll 1$

$$\Gamma = \frac{\psi}{V_i} k^2 (V_d^2 - C_s^2)$$

$$\omega_r = k \cdot V_d$$

THRESHOLD: $\omega_r/k = C_s$

$f = 49.92 \text{ MHz}$

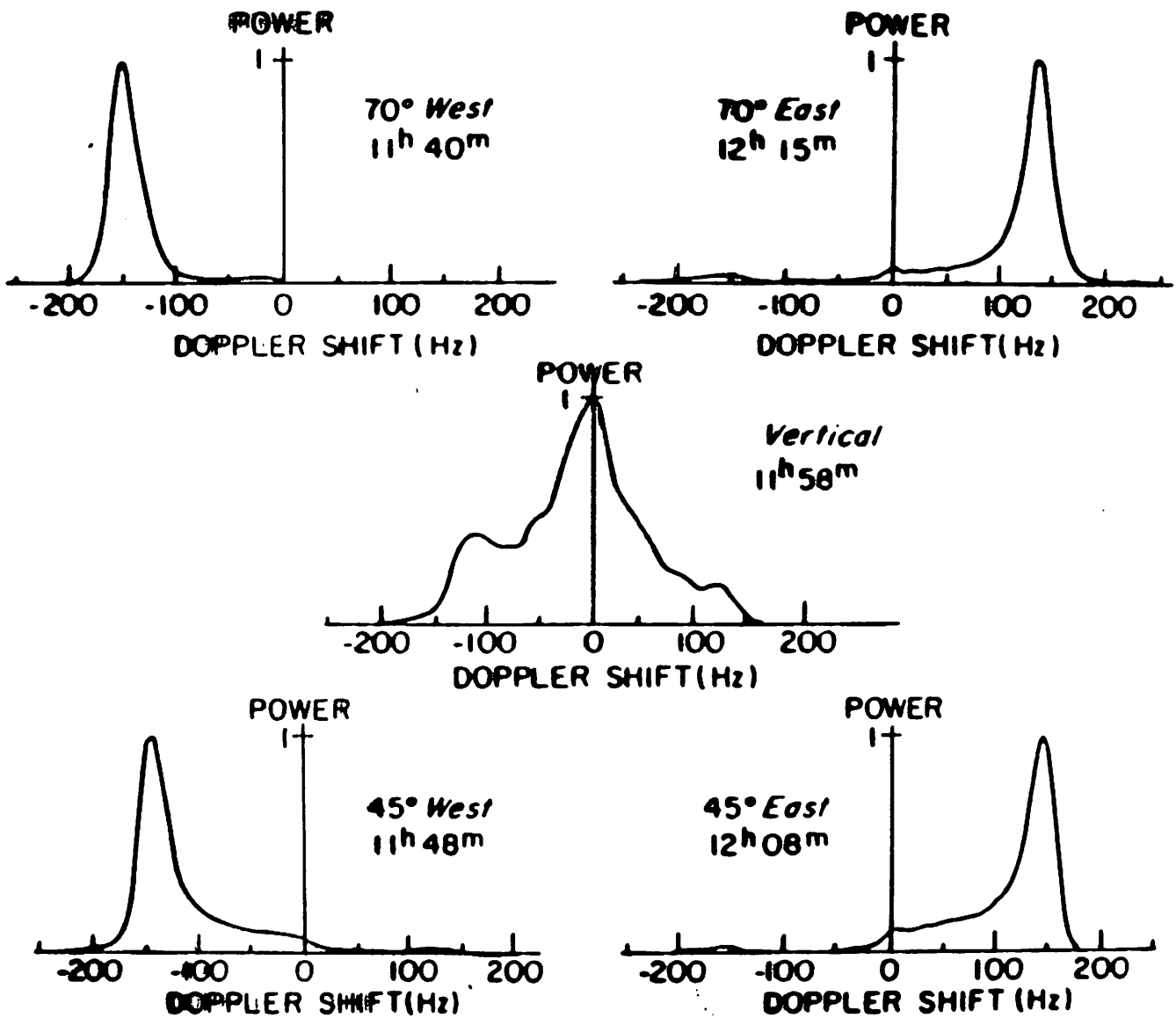
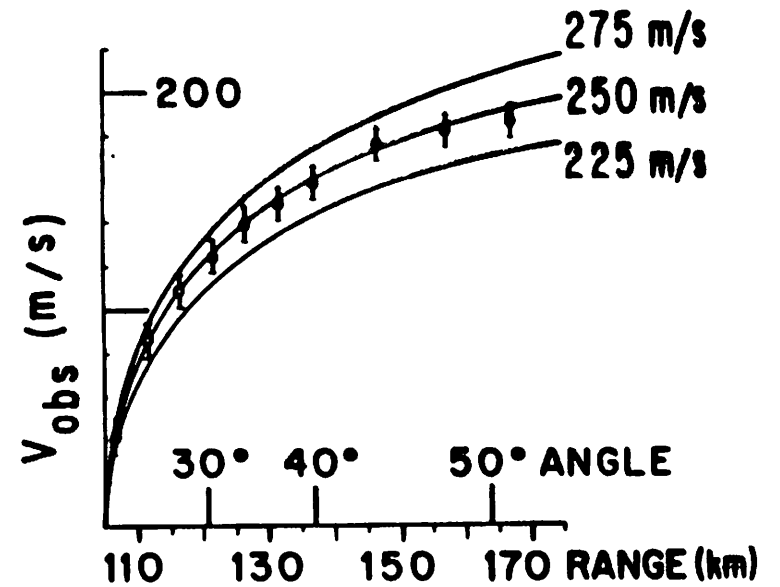
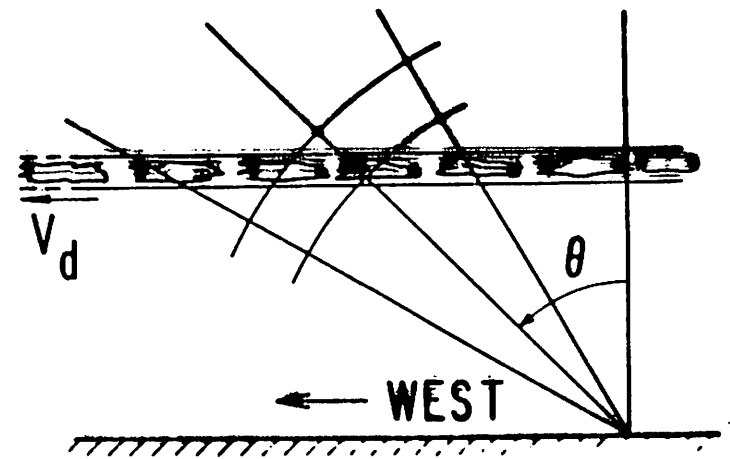
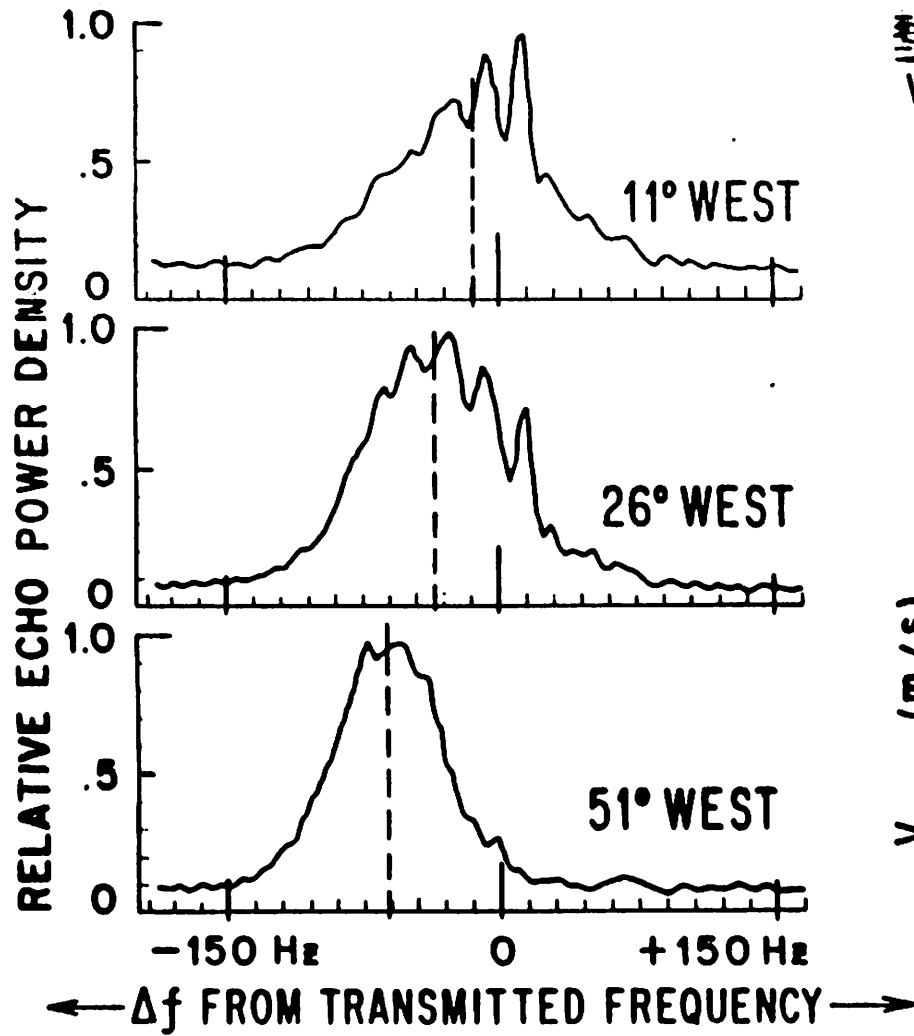
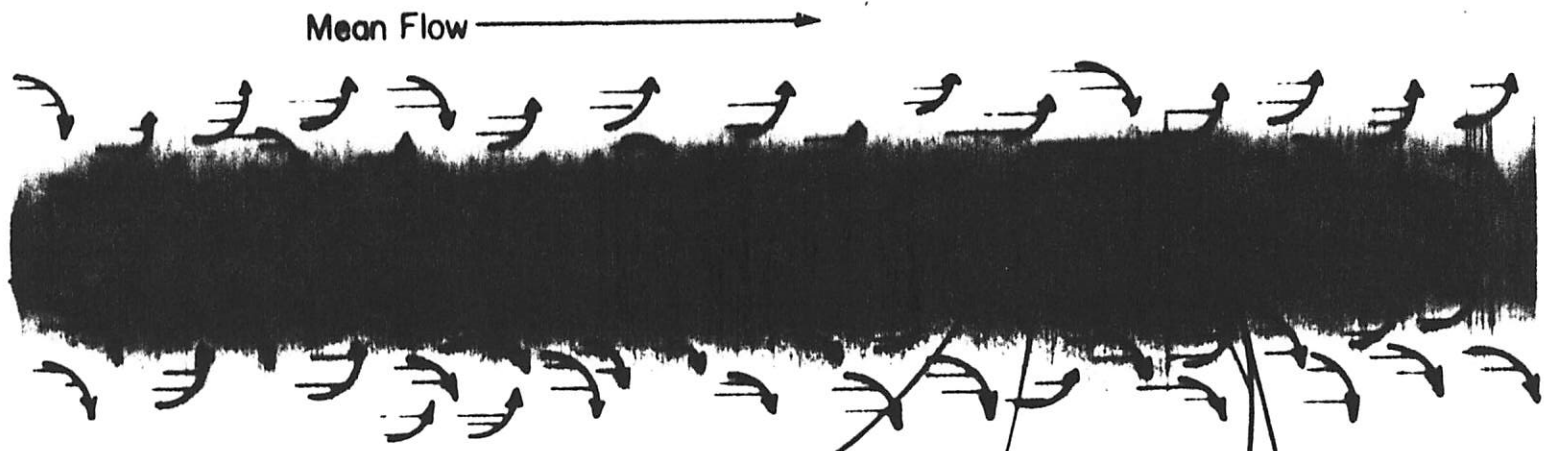


Fig. 4.20. Series of Doppler spectra from the equatorial electrojet irregularities at different elevation angles obtained at Jicamarca during a period of relatively strong scattering. The spectra are normalized to a fixed peak value. [After Cohen and Bowles (1967). Reproduced with permission of the American Geophysical Union.]

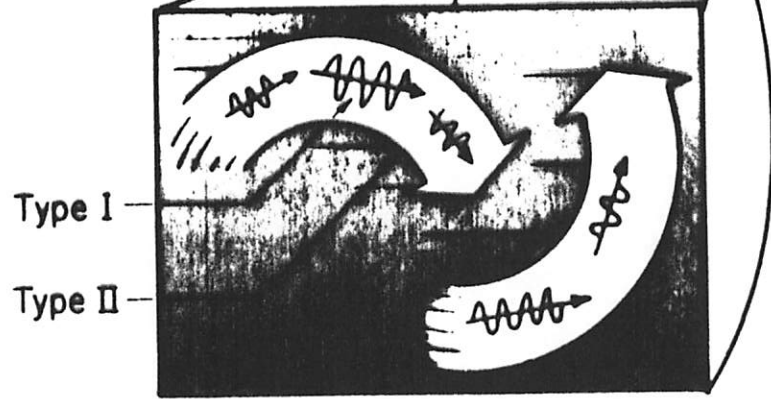
JICAMARCA, PERU

16 NOV 1966, 17^h 53^m (75° W)

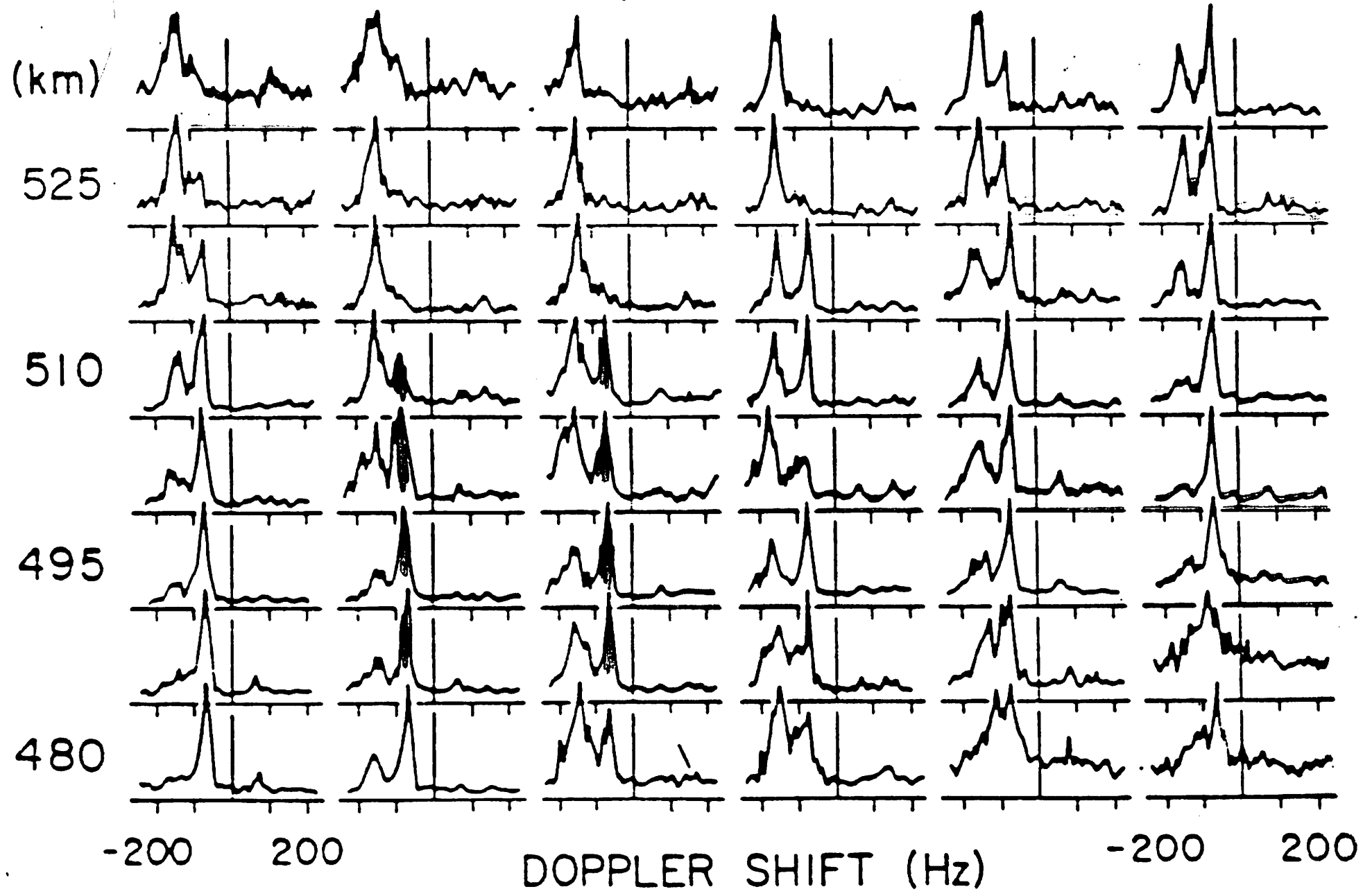




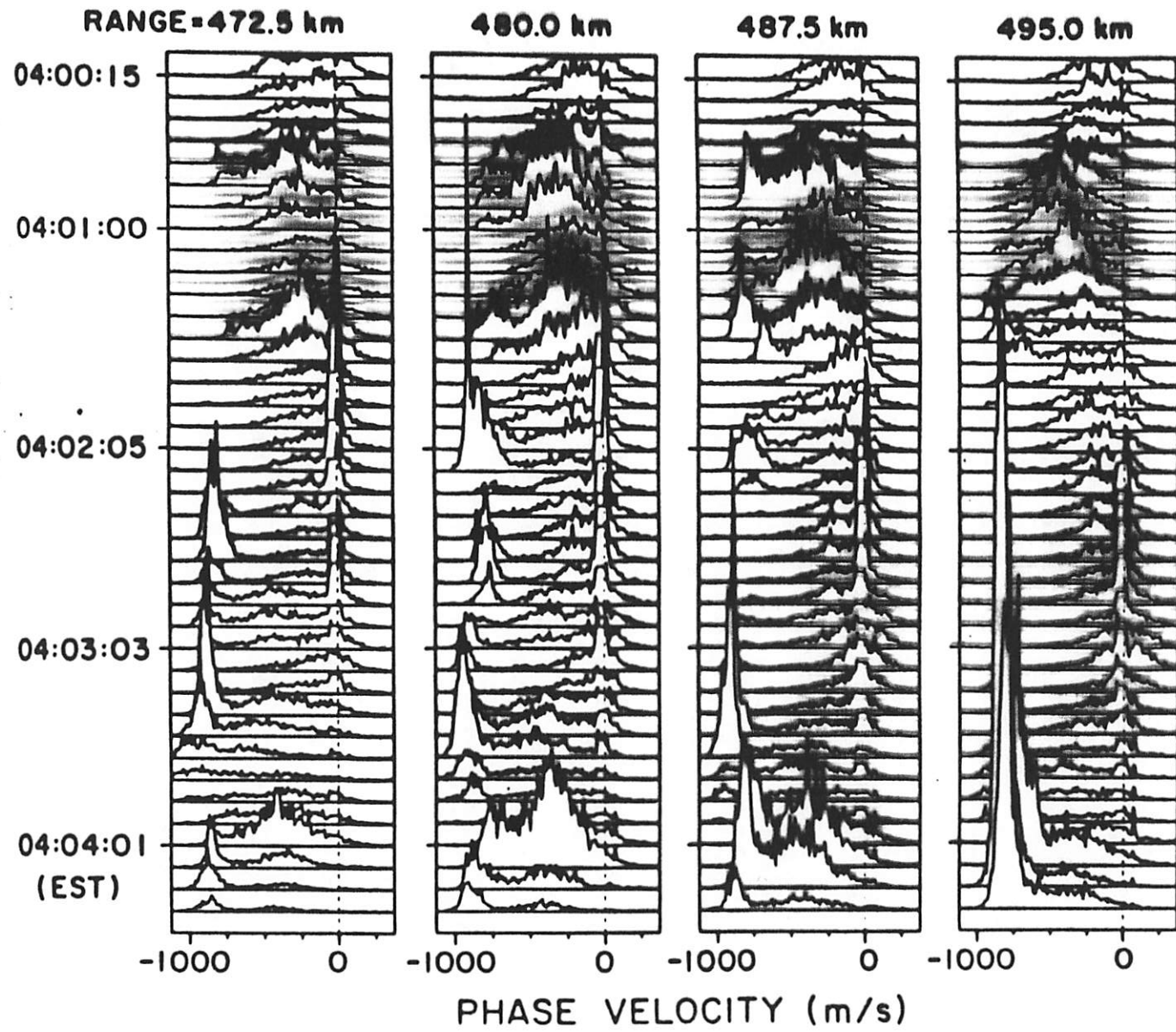
Electrojet Region



ITHACA, N.Y. APRIL 1, 1976 02:04:15-02:04:28 LT



ITHACA, NY SEPTEMBER 6, 1982



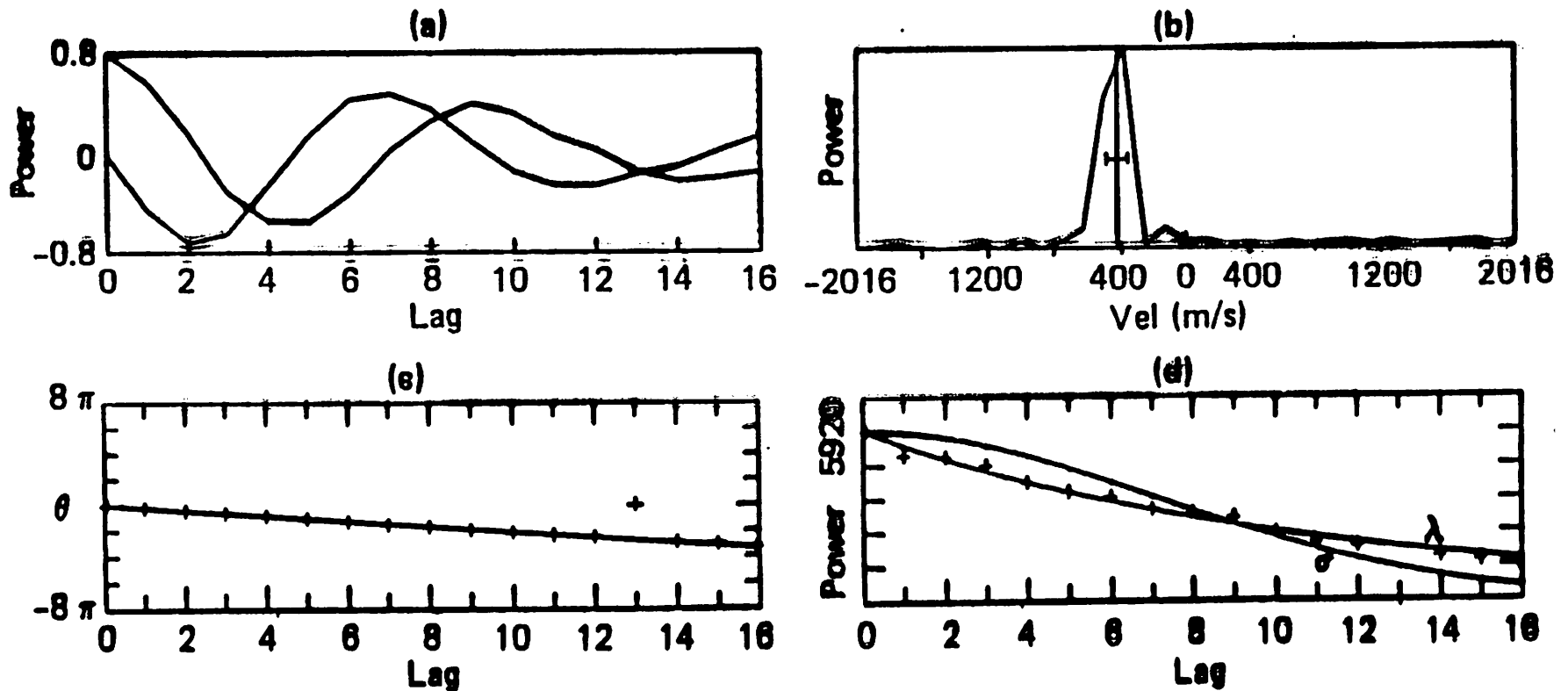
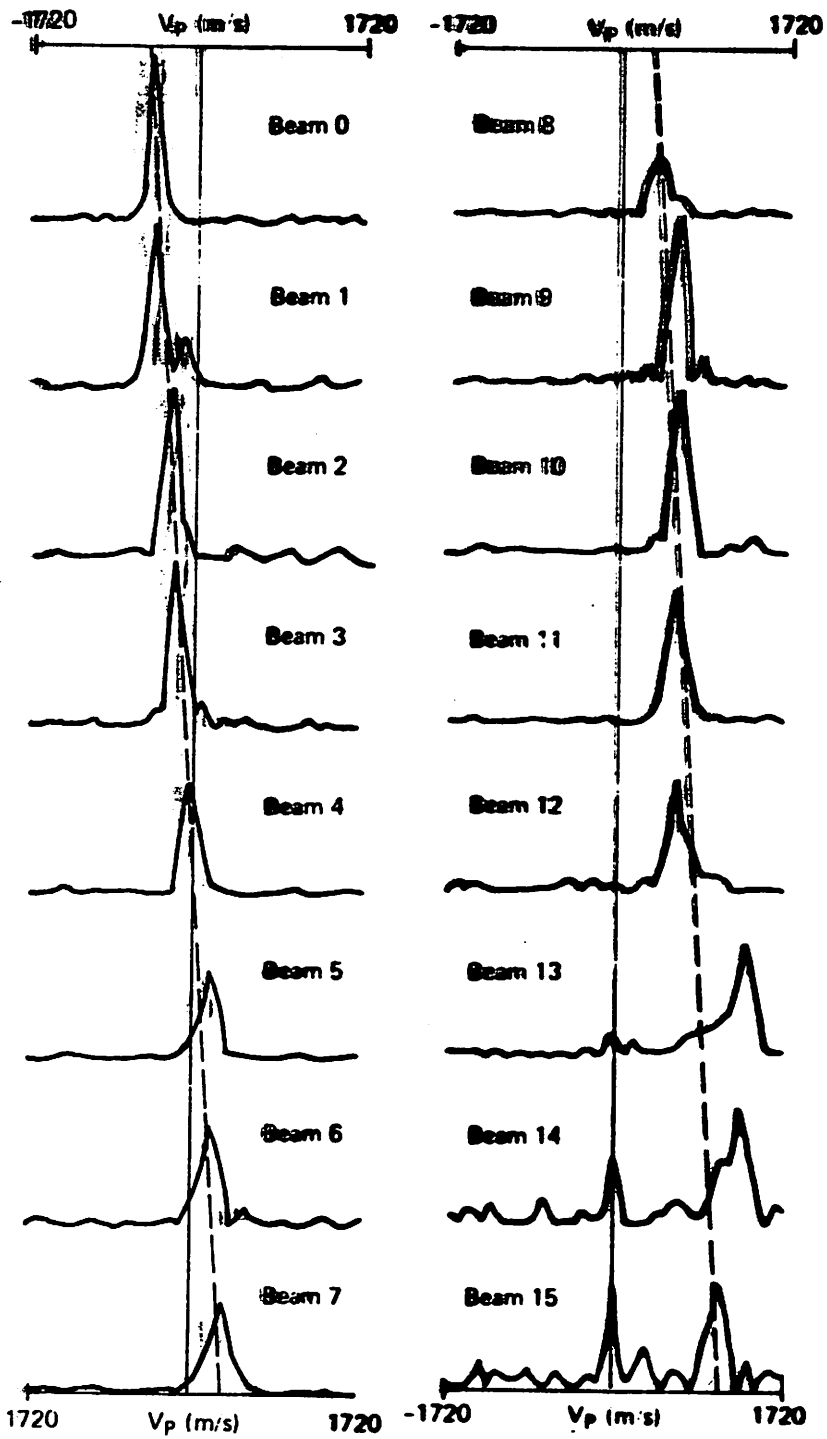
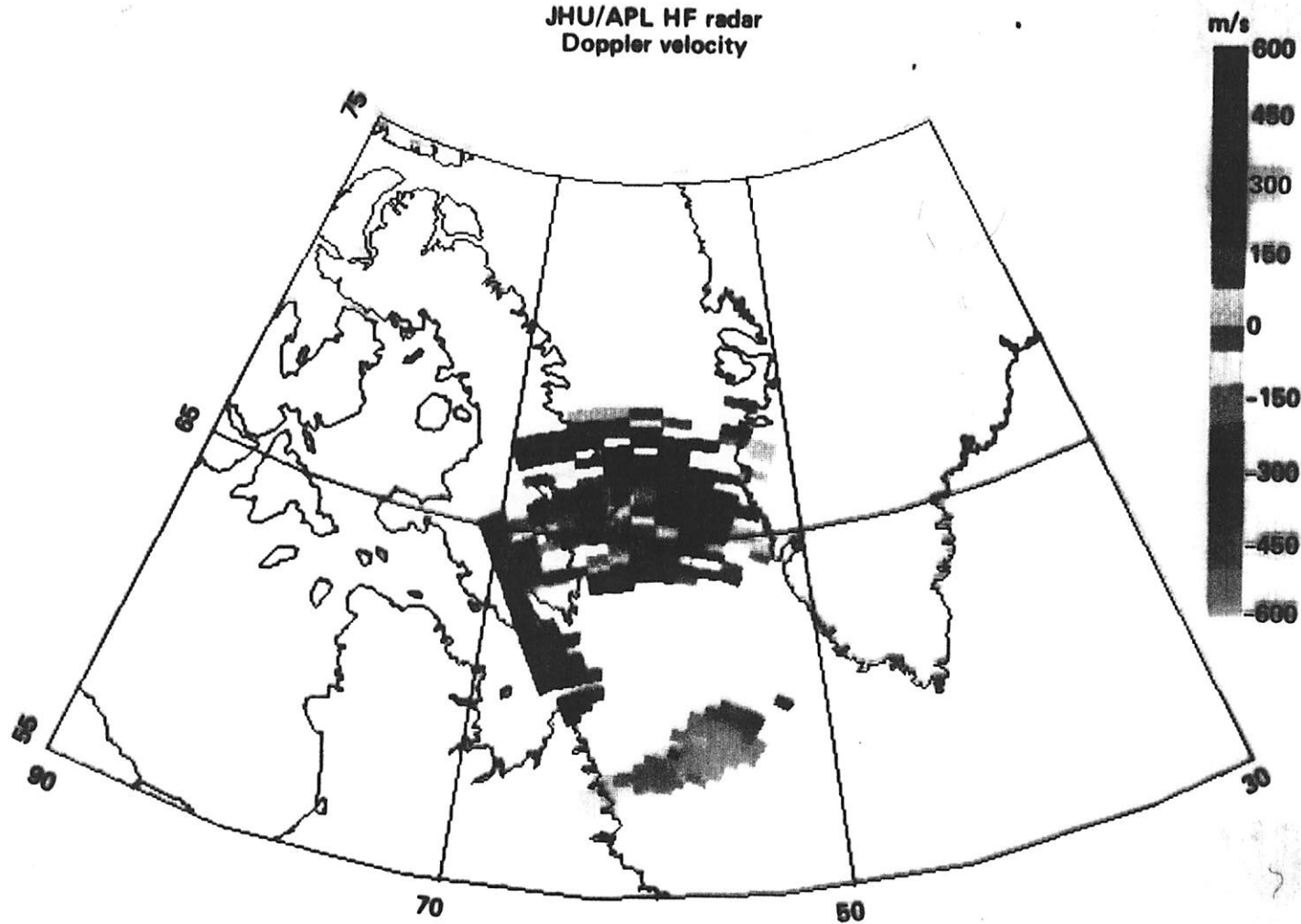


Figure 5. Stages in the processing of a sample of HF radar data. (a) Real and imaginary parts of an ACF. (b) Doppler spectrum obtained by Fourier transformation of the ACF. (c) Phase angle as a function of lag and its linear least-squares fit. (d) Power variation of the ACF as a function of lag and fits based on exponential λ and exponential $\sigma = \lambda^2$ as explained in the text. [Reprinted from Villain et al., 1987]



12 October 1983 2001:32 - 2002:52 UT
 Frequency = 4.5 MHz Range = 1100 km

JHU/APL HF radar
Doppler velocity



Frequency: 12.3 MHz

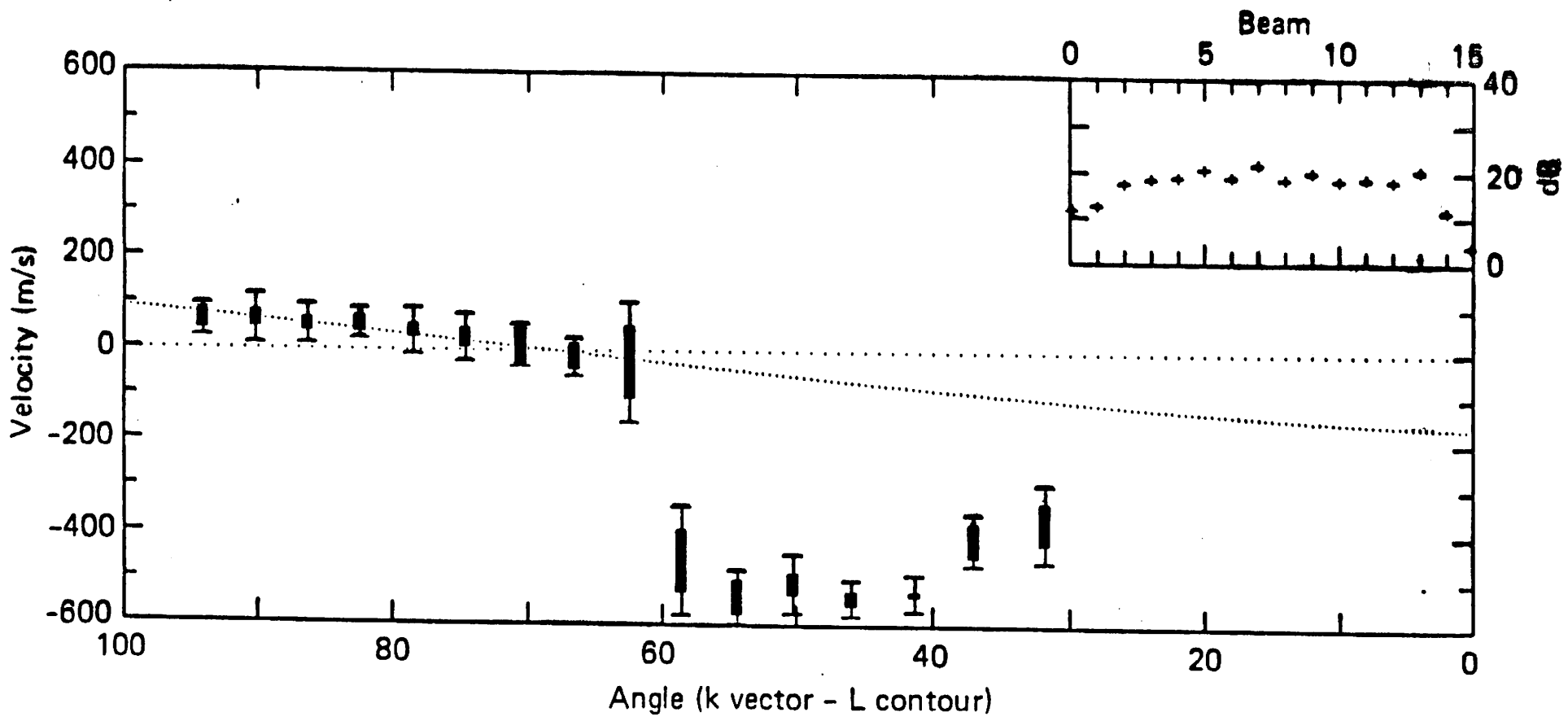
Date: 9-10-86

Time: 05:51:35

Date: 09-10-86 Start time: 05:51:35 Freq: 12.4 MHz

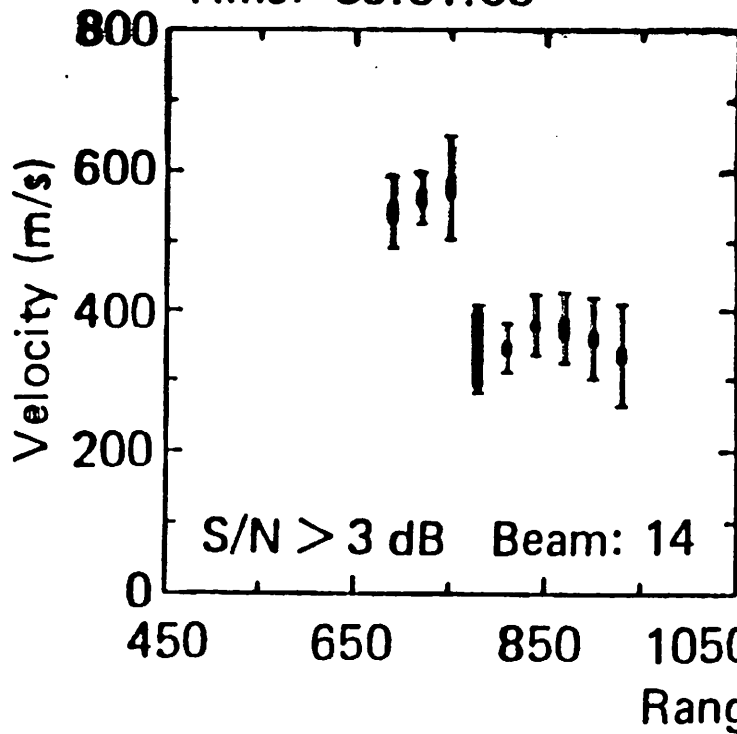
Beam: 00 01 02 03 04 05 06 07 08 09 10 11 12 13 14 15

Range: 09 09 09 09 09 09 10 10 11 11 12 13 14 16 17 20



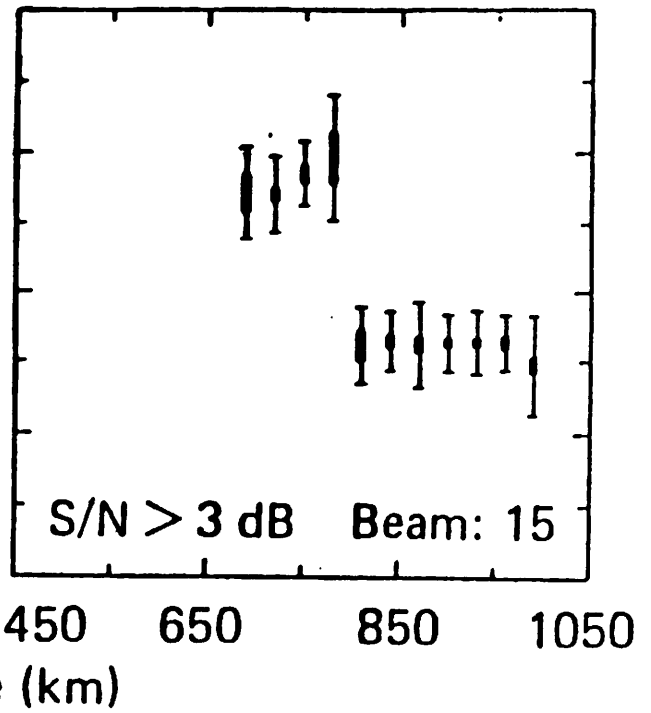
Date: 09-10-86

Time: 06:01:00



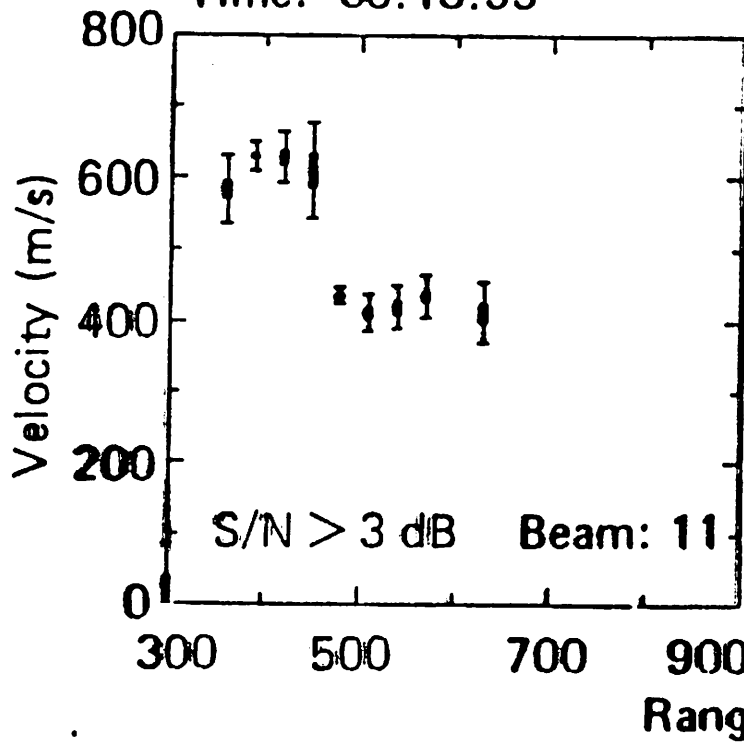
Date: 09-10-86

Time: 06:01:05



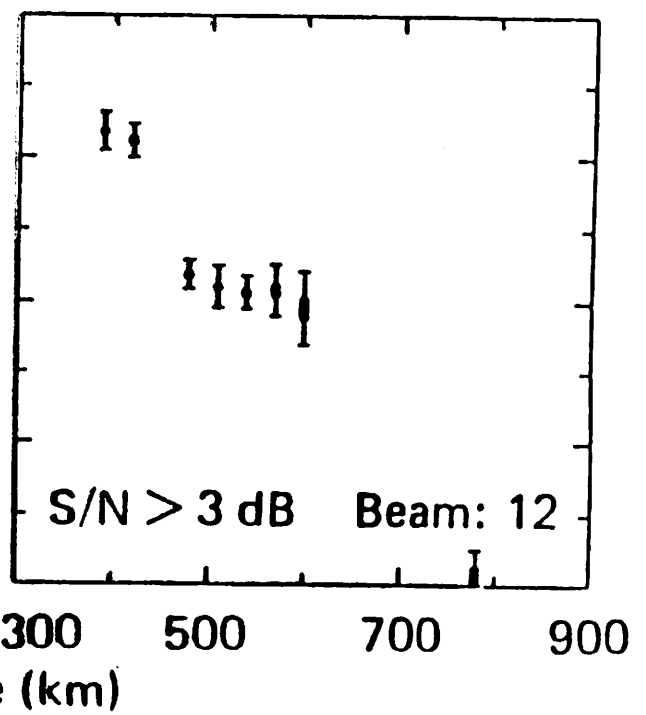
Date: 10-04-85

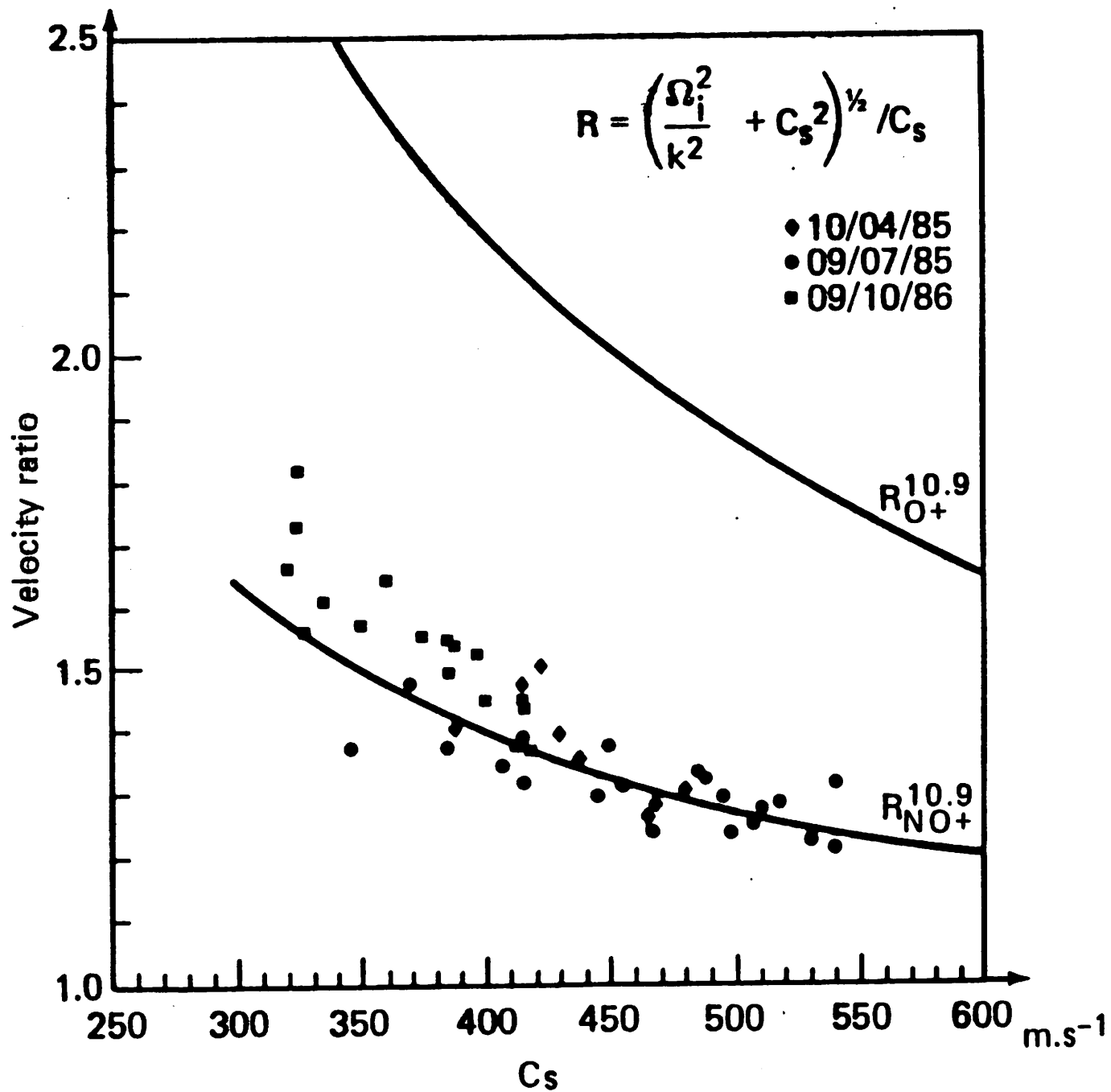
Time: 06:18:55



Date: 10-04-85

Time: 06:19:00







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DAYSIDE CONVECTION

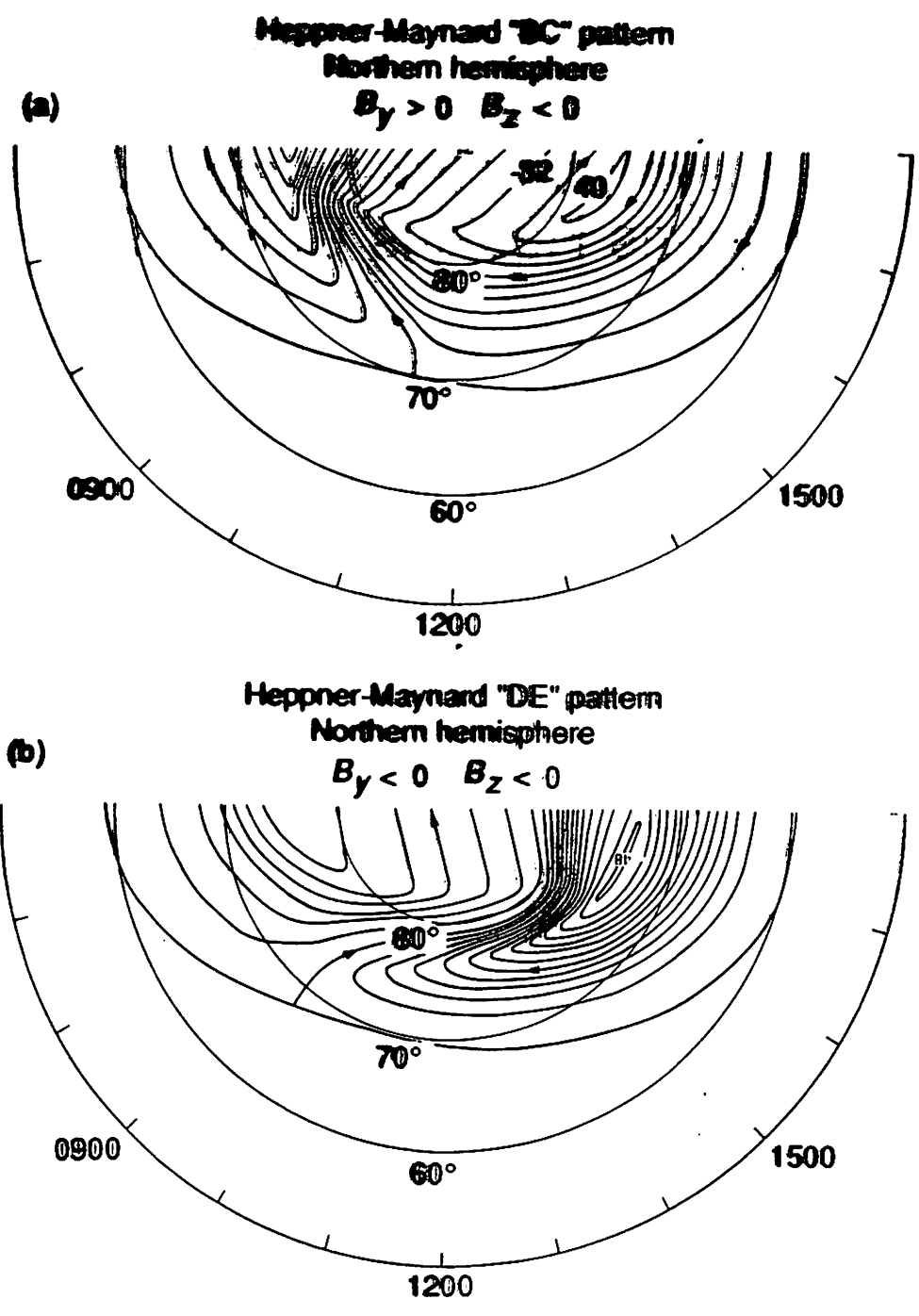
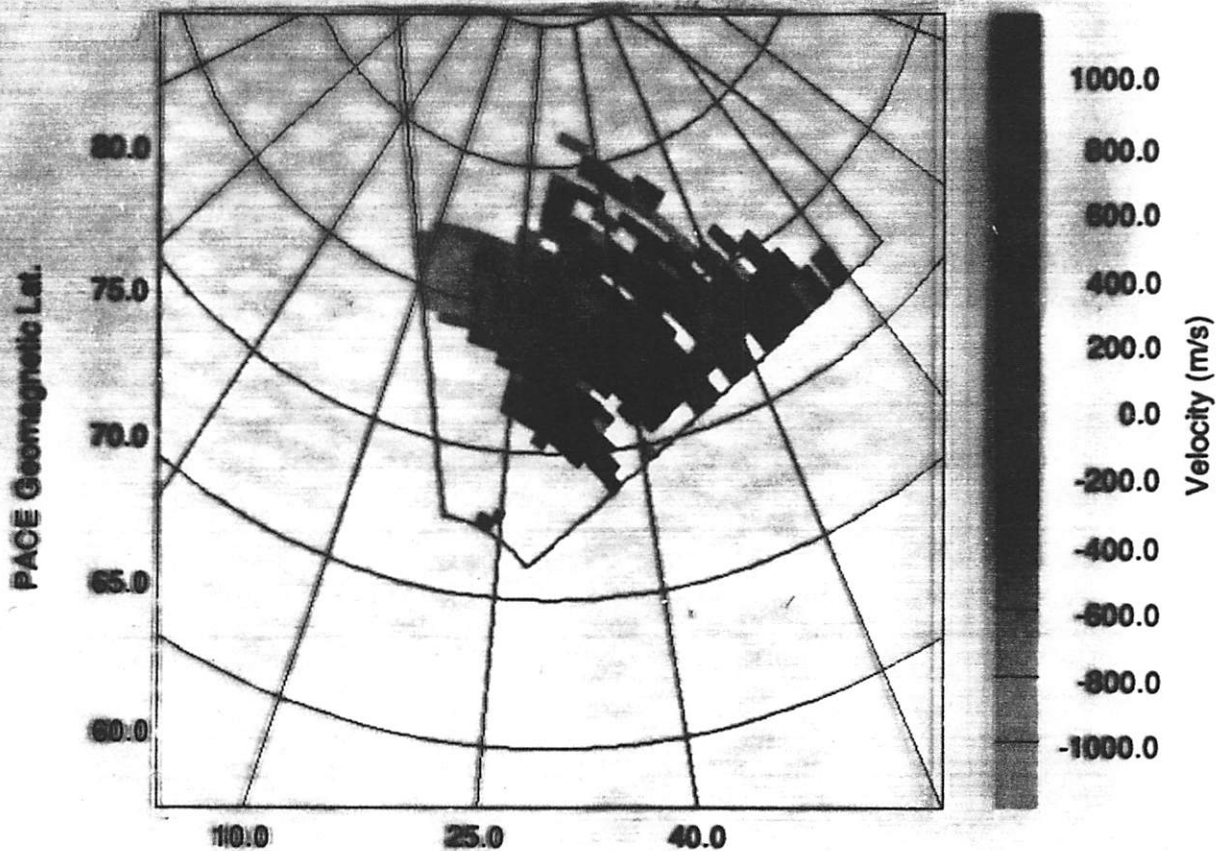
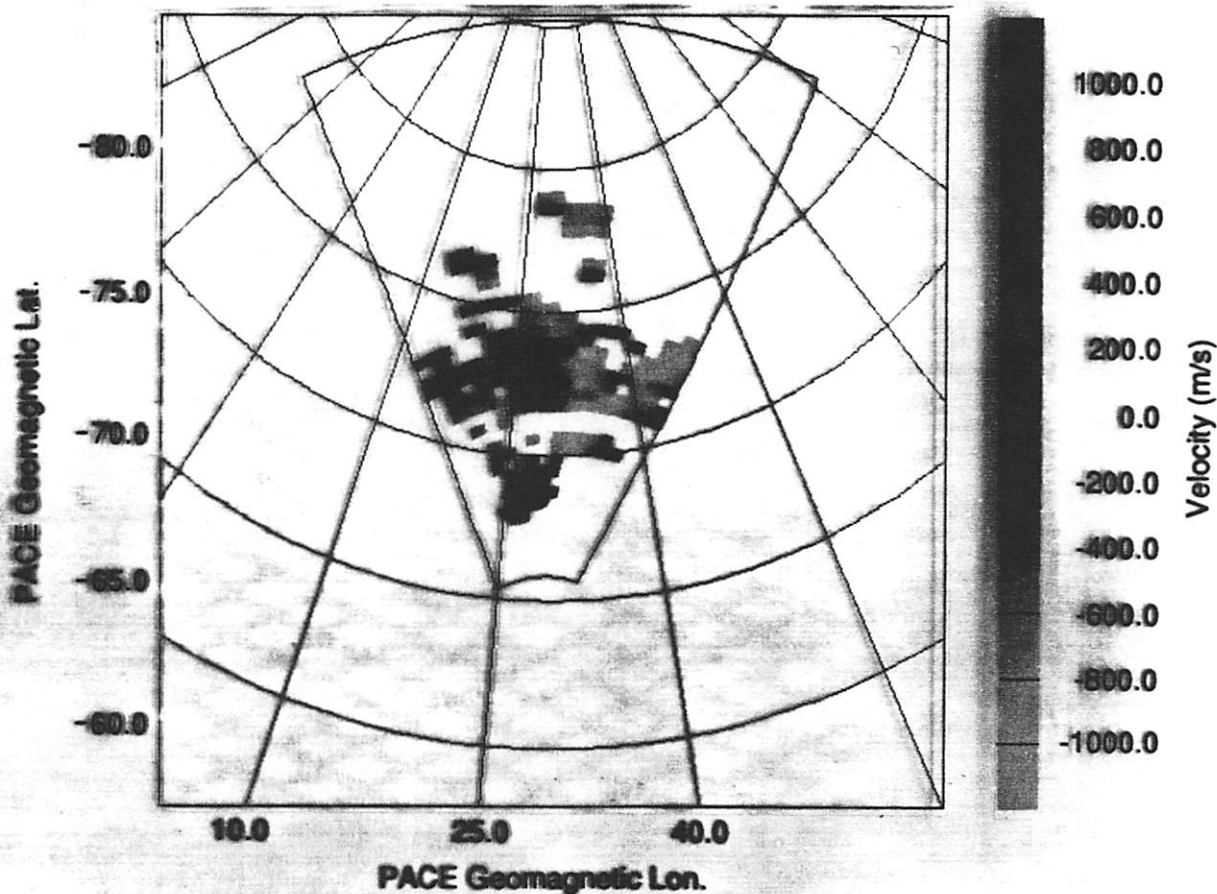


Fig. (1)

GOOSE BAY
Time: 15:31:17, Freq: 12.4 MHz.

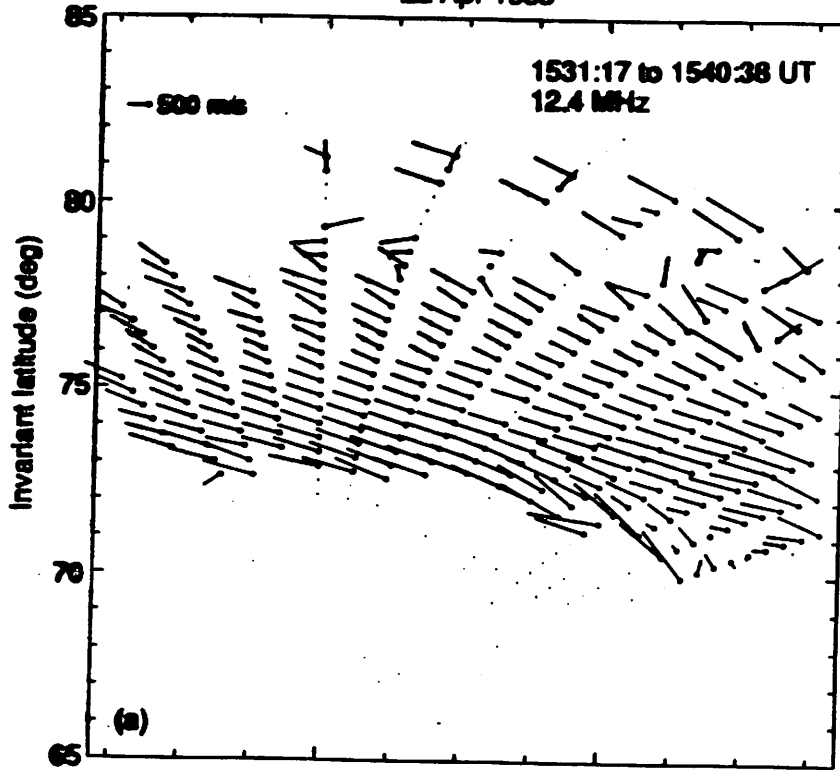


HALLEY
Time: 15:31:52, Freq: 12.2 MHz.



PACE Radar Comparison Plot 22/04/1988.

Goose Bay HF radar
22 Apr 1988



Halley HF radar
22 Apr 1988

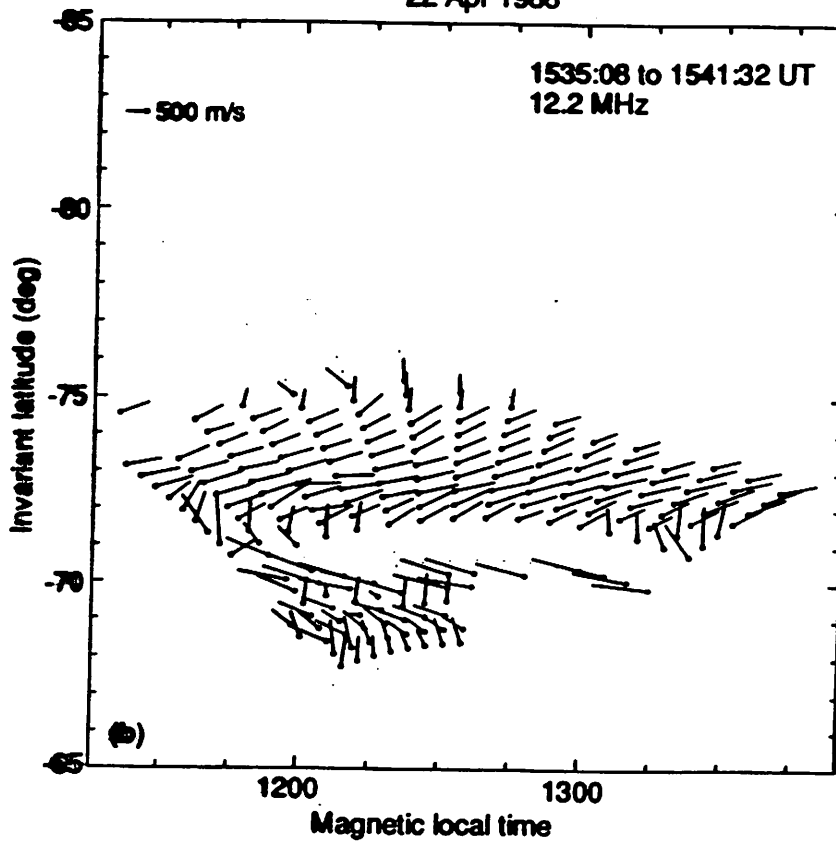


Fig. 6

IMP-8 Magnetometer

22 Apr 1988

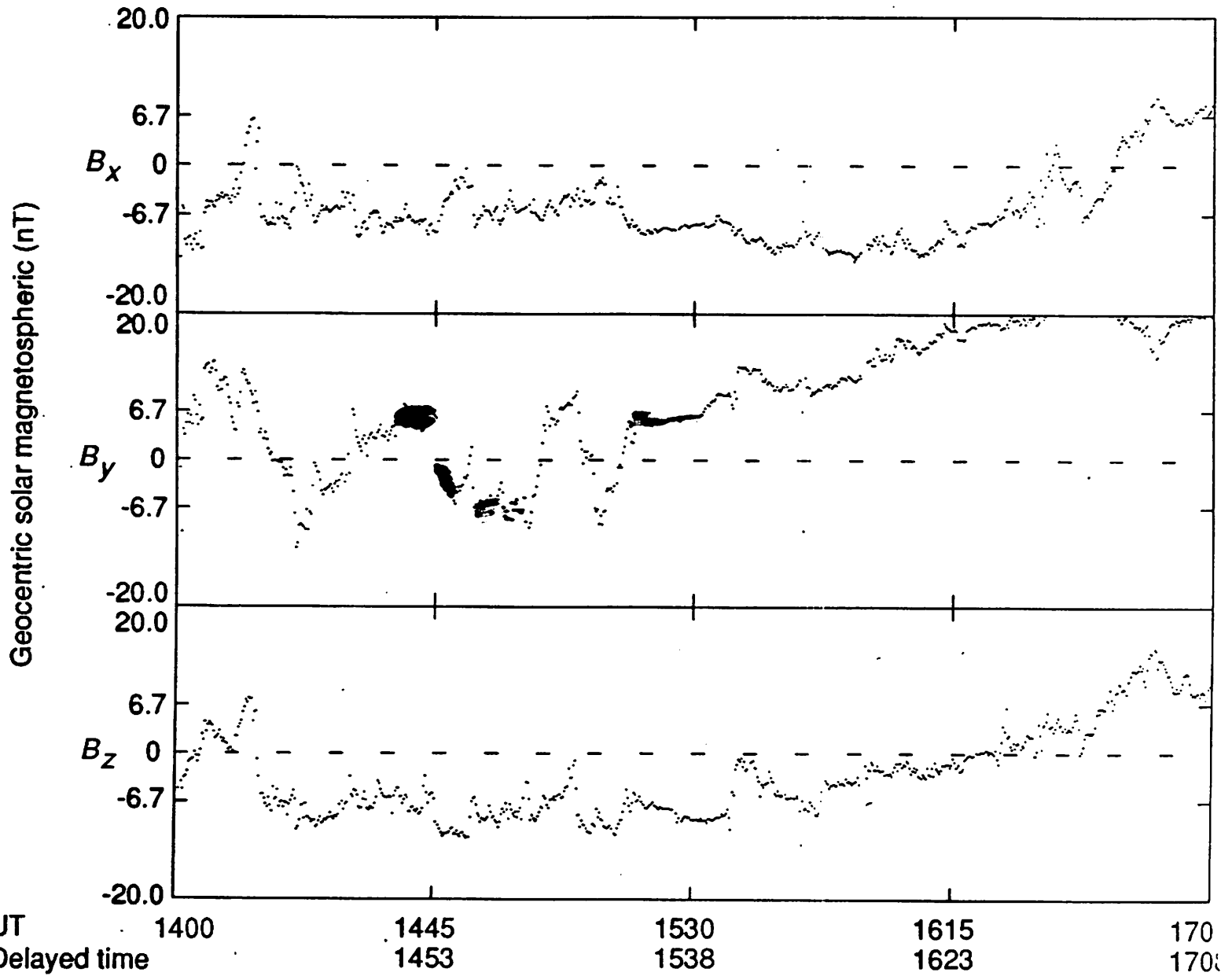
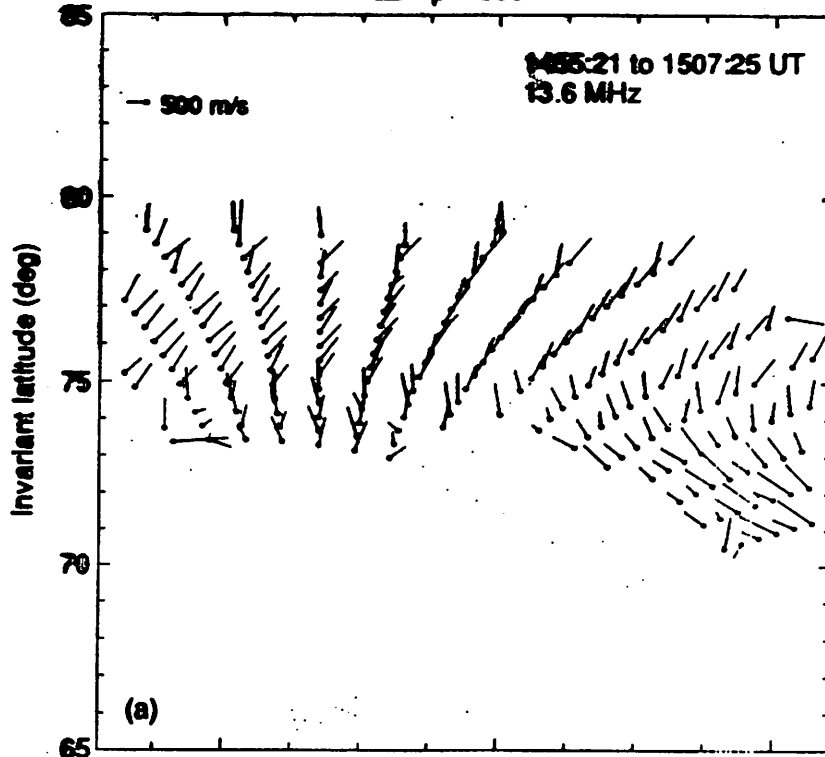


Fig. 3

Goose Bay HF radar
22 Apr 1988



Halley HF radar
22 Apr 1988

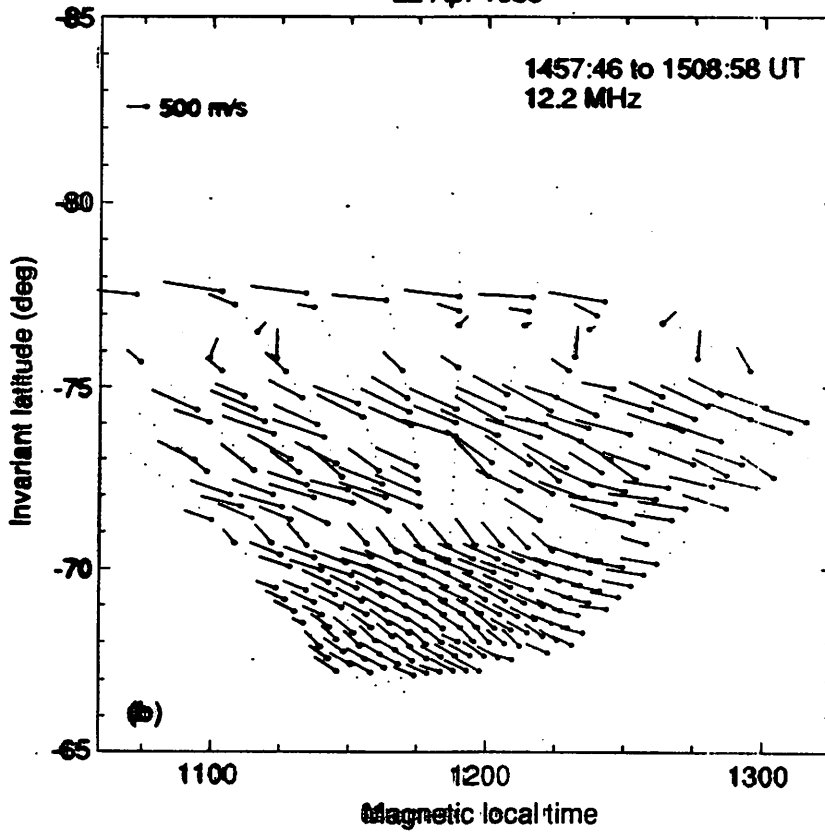


Fig. 7

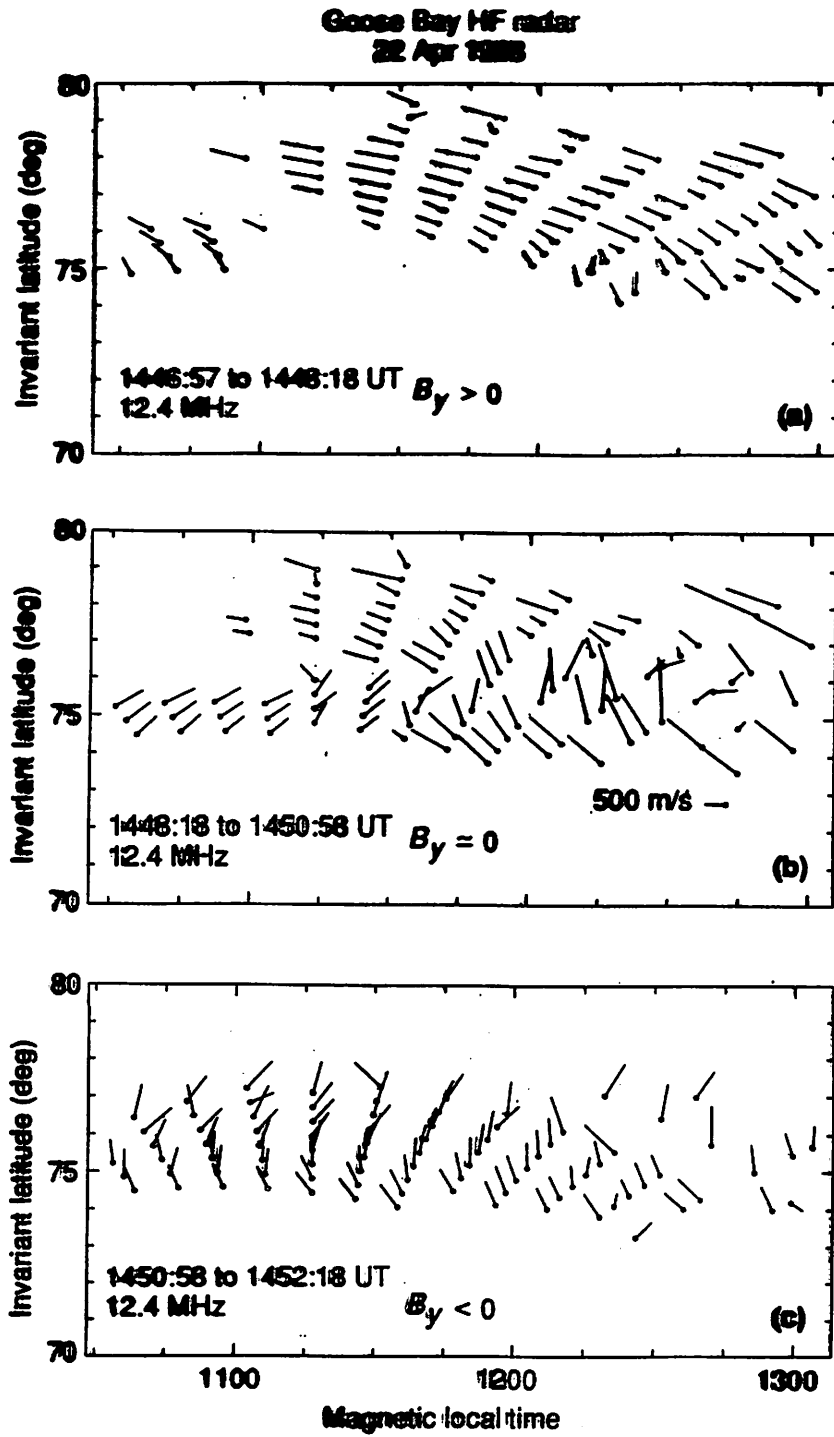


Fig. 9

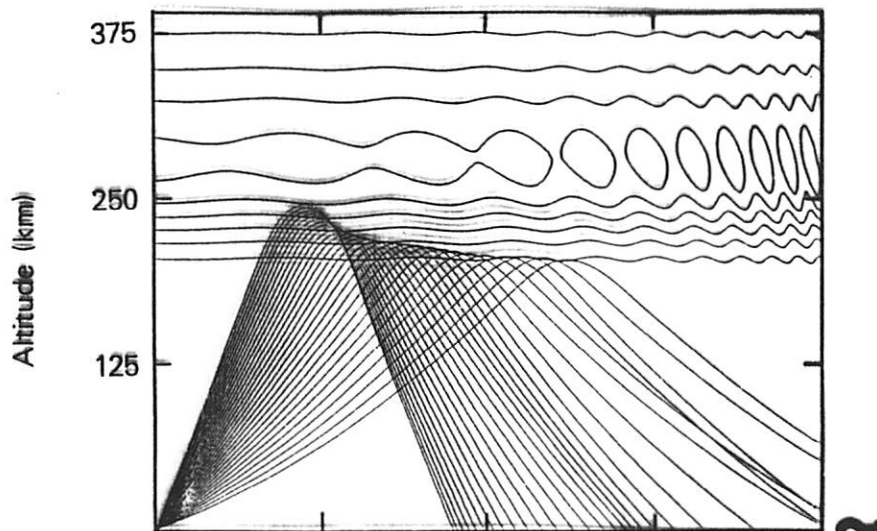


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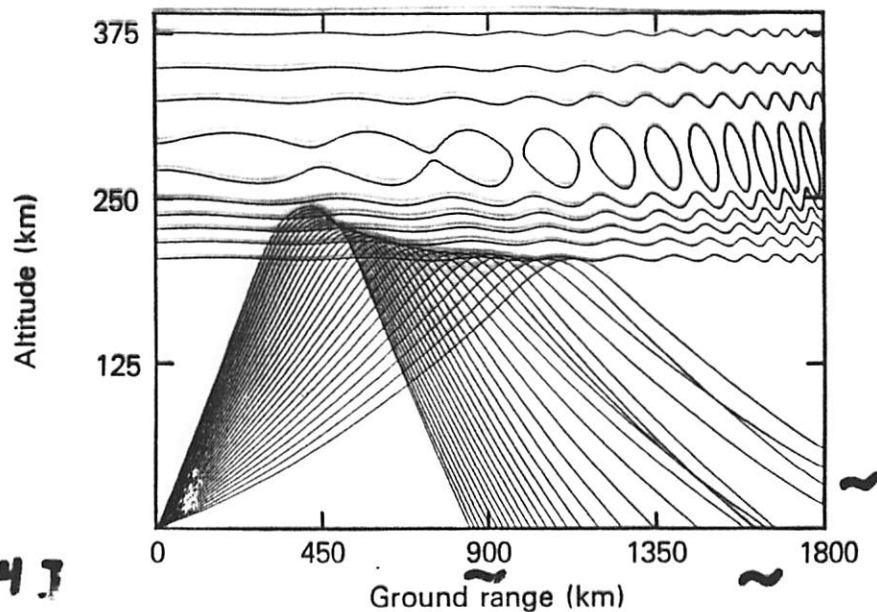
GRAVITY WAVES

MODELLING GRAVITY WAVE EFFECTS:

HF RAY TRACING



~ focusing

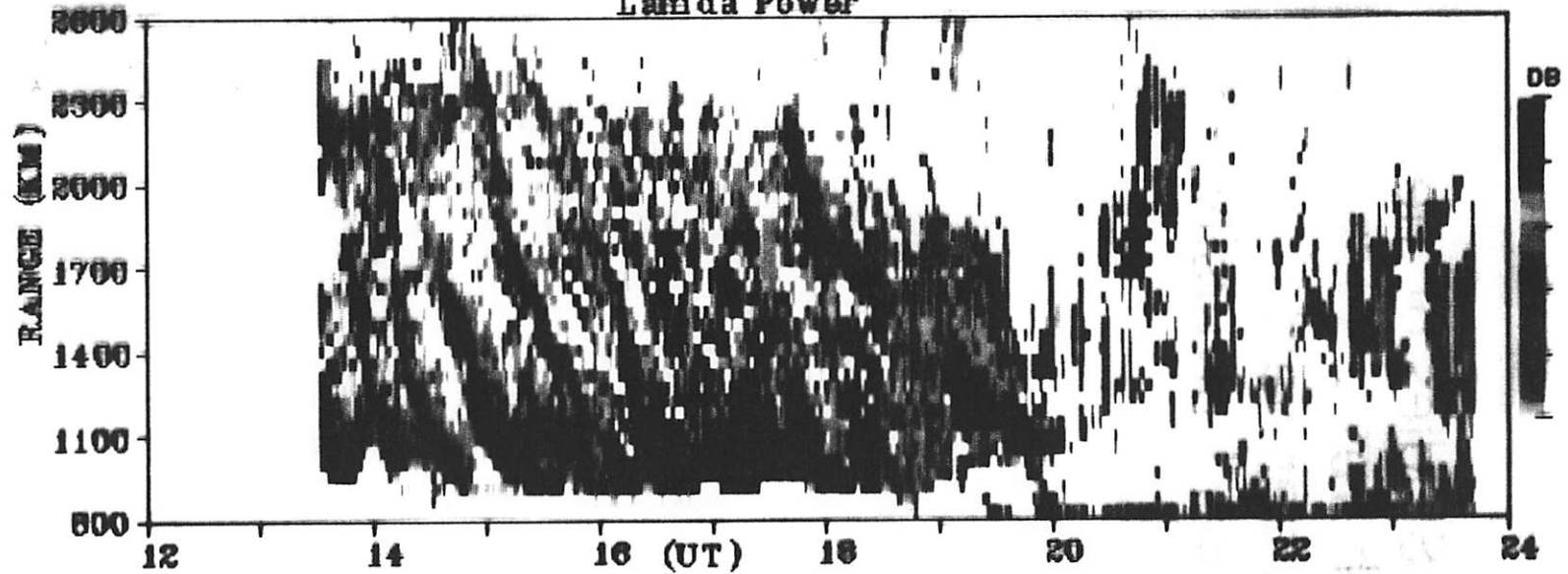


+ 15 MIN

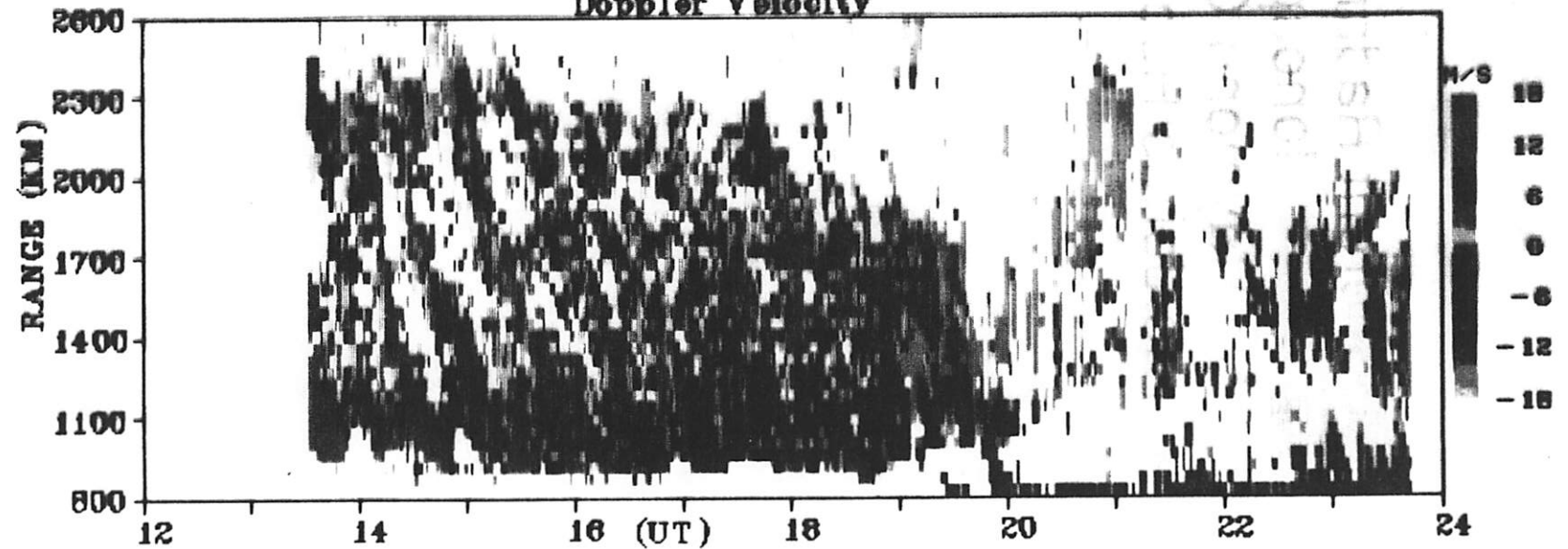
TID due to
Earth-Reflected
Gravity Wave
[Francis, JGR, 1974]

JHU/APL HF-RADAR

Lambda Power



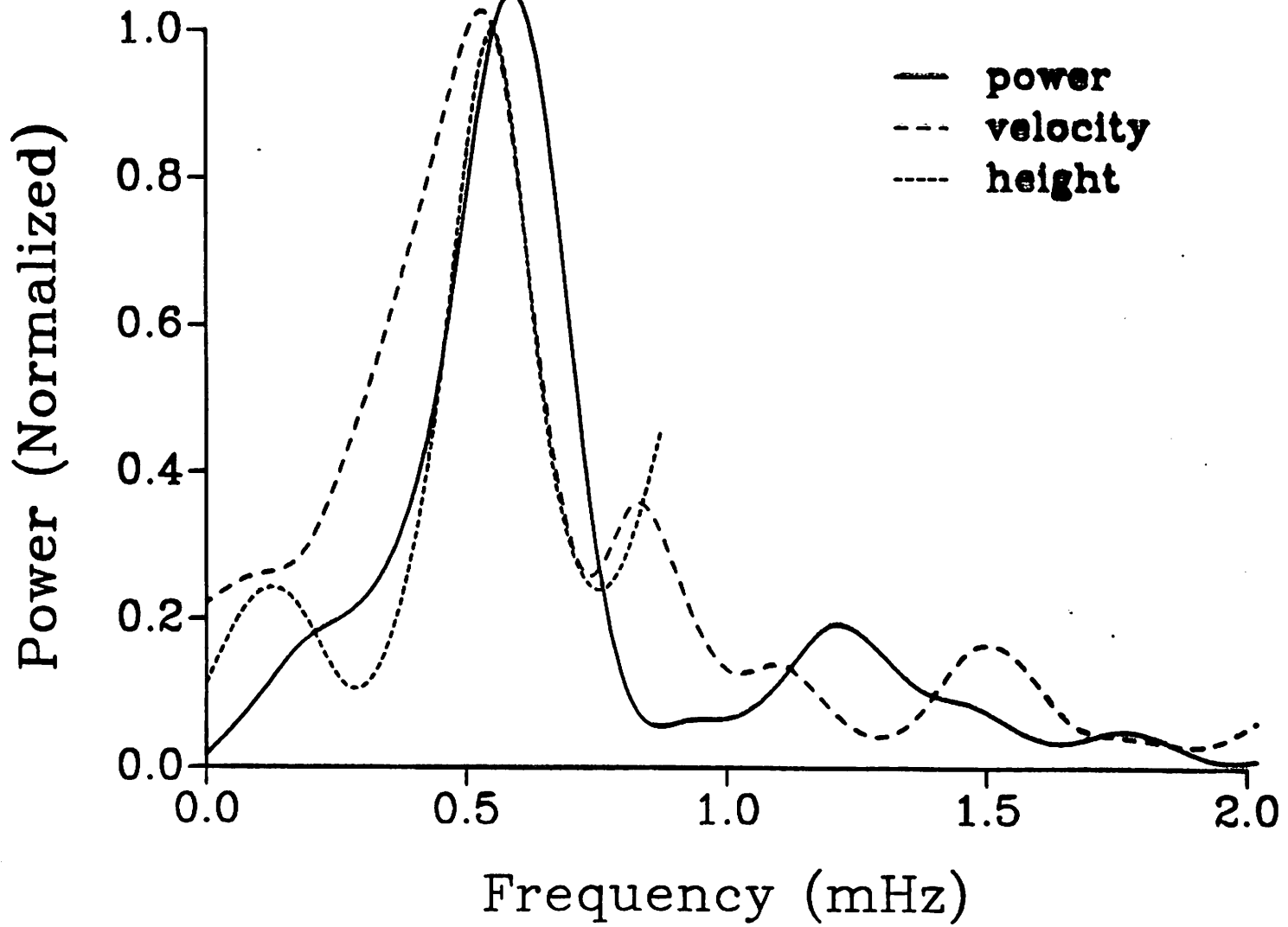
Doppler Velocity

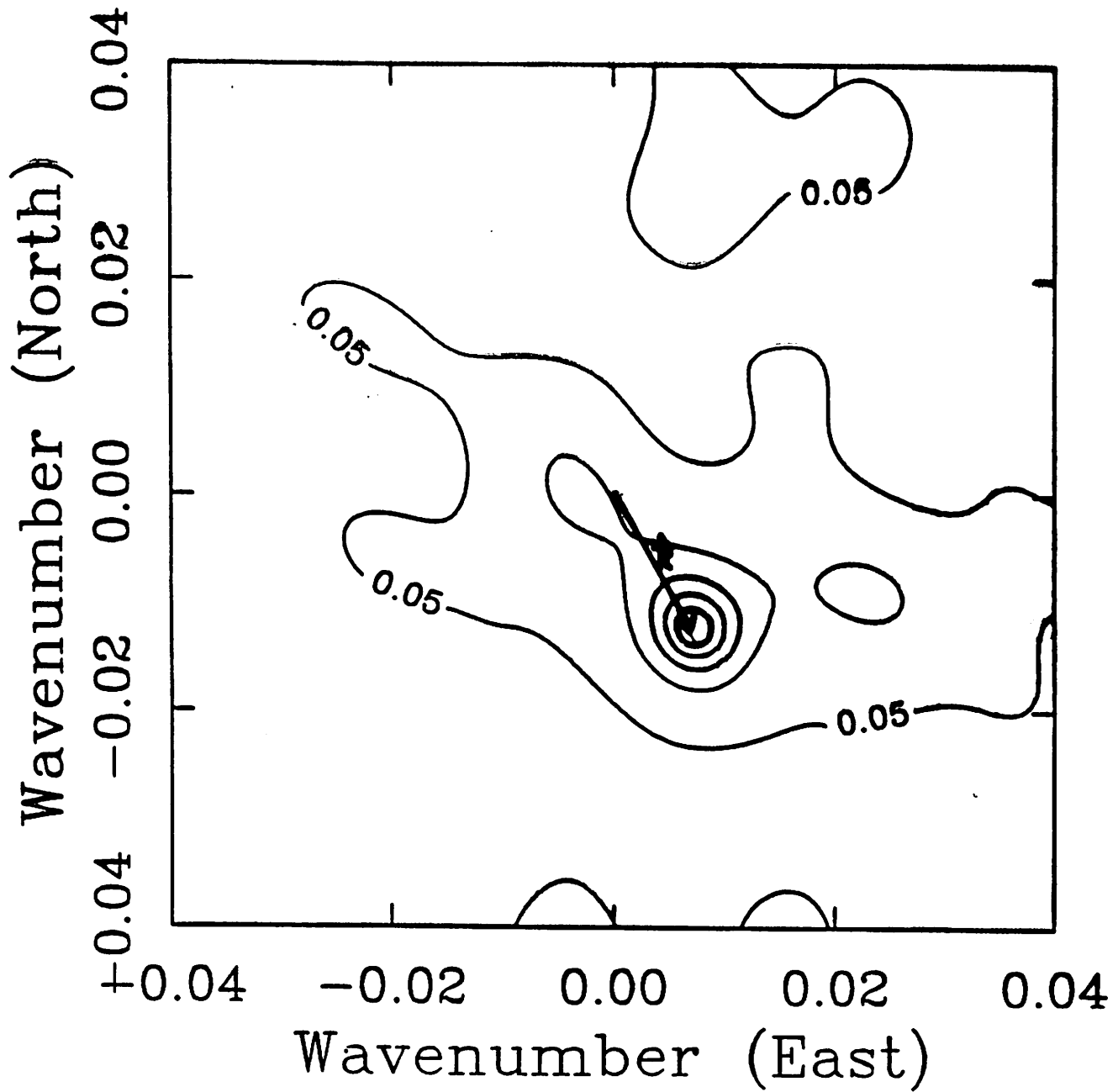


BEAM : 8

FREQUENCY : 14.6 MHz

DATE : 11:25:1988



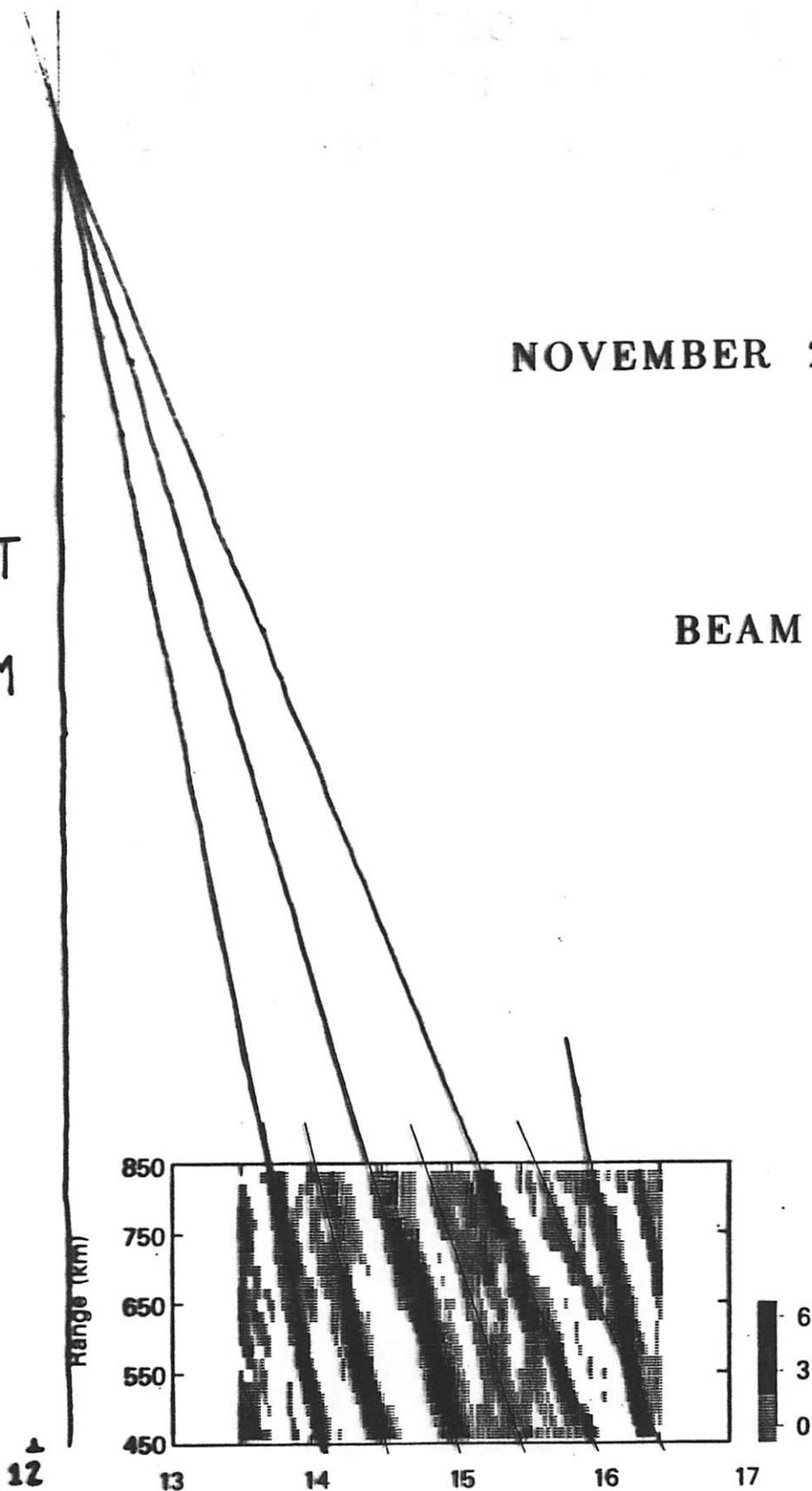


MODELLING GRAVITY WAVE PROPAGATION:
THEORETICAL AND EXPERIMENTAL PHASE MAPS

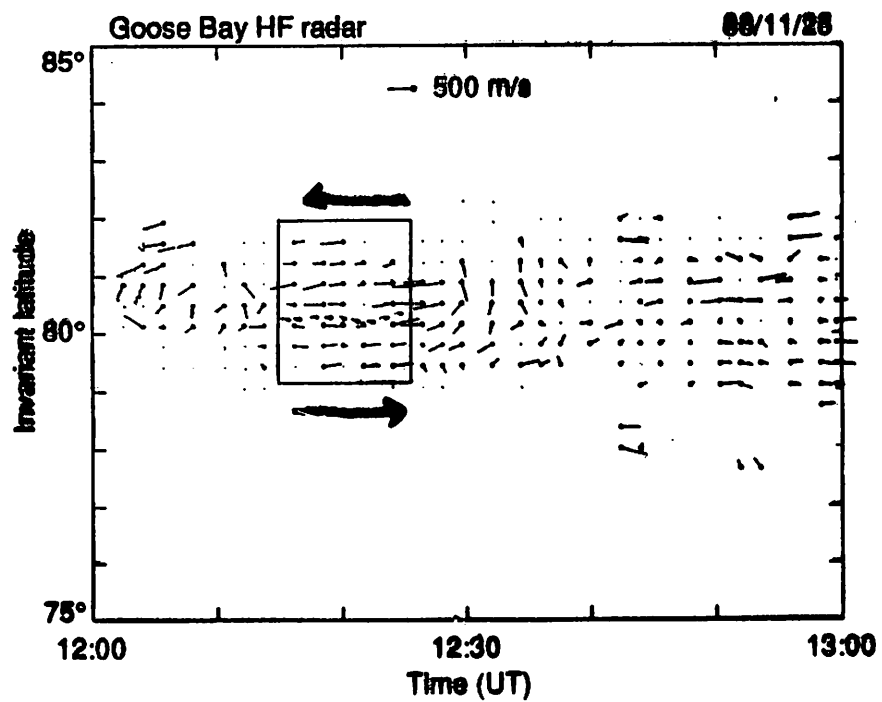
NOVEMBER 25, 1988

ONSET:
~ 1211 UT
~ 2100 KM

BEAM 8



GENERATION OF A GRAVITY WAVE:
LATITUDINAL PROFILE OF CONVECTION VELOCITY



NOVEMBER 25, 1988 1200 — 1300 UT

SuperDARN

*Global Imaging of
Ionospheric dynamics*



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What Is SuperDARN?

SuperDARN is an internationally-funded network of HF radars capable of continuous, global-scale observations of the high-latitude ionosphere.



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What is the basis of the SuperDARN concept?

Bidirectional common volume observations with pairs of HF radars to provide vector determinations of plasma convection and the ionospheric electric field

**Spatial resolution of each vector
~50 km x ~50 km**

**Total spatial coverage
~40 million km²**

Observations will cover the auroral zone and polar cap

**$65^\circ < \Lambda < 85^\circ$
~14 Hours of Magnetic Local Time**

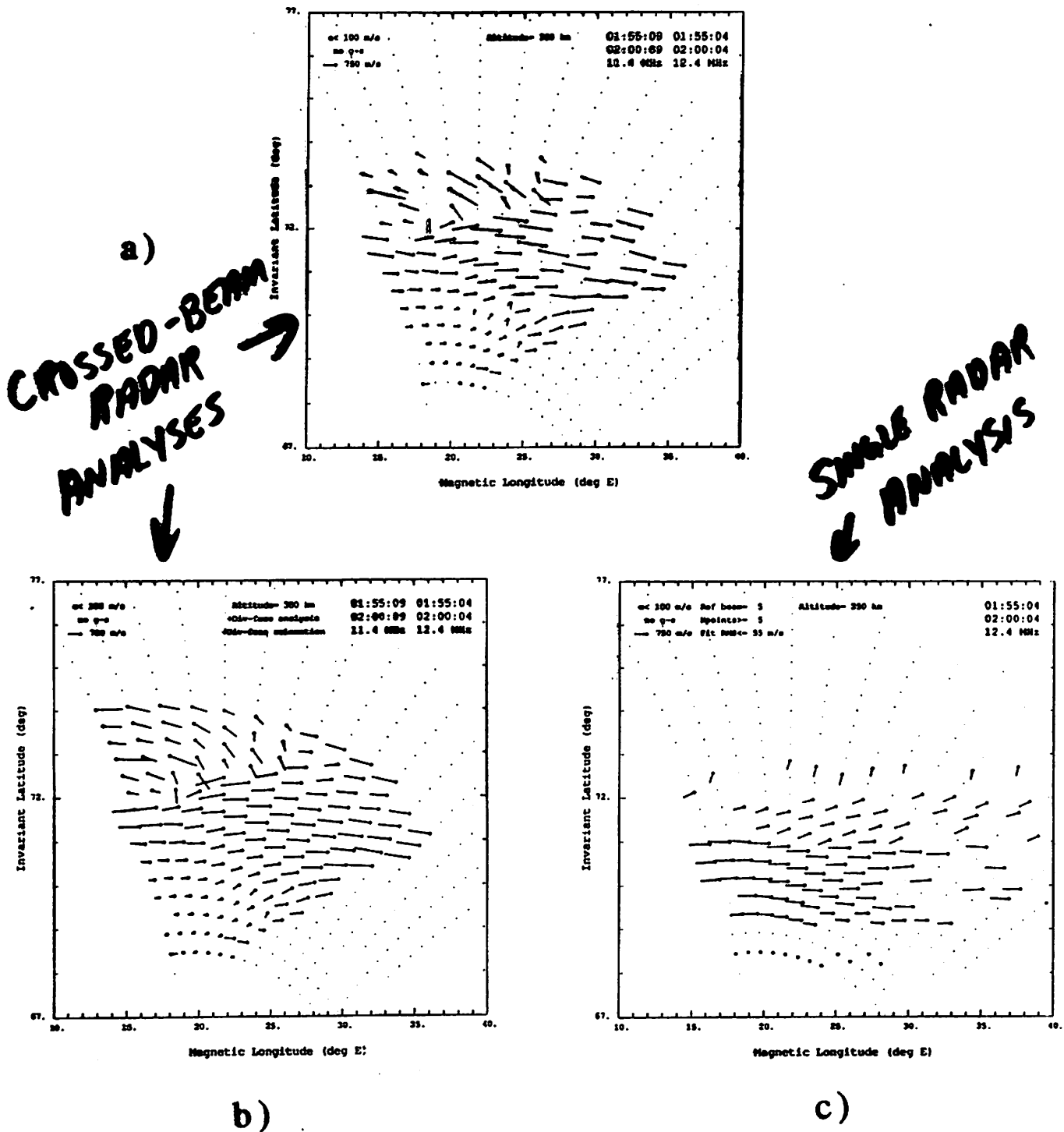


Figure 7 a) Map of convection velocity for the period December 6, 1989, 01:55 - 02:00 UT obtained by merging line of sight velocity data from the Goose Bay and Schefferville HF radars. b) Same as (a), except that the region of vector solution has been extended by applying the condition that the flow be free of divergence to Goose Bay line of sight data, followed by a relaxation of the overall pattern to a divergence-free state. c) Map of convection velocity for the period of (a) that is obtained by fitting the Goose Bay line of sight velocity data on the basis of an assumed L-shell invariance in the flow. Moderate fitting criteria have been applied.



SuperDARN (Northern Hemisphere) Division of Responsibilities

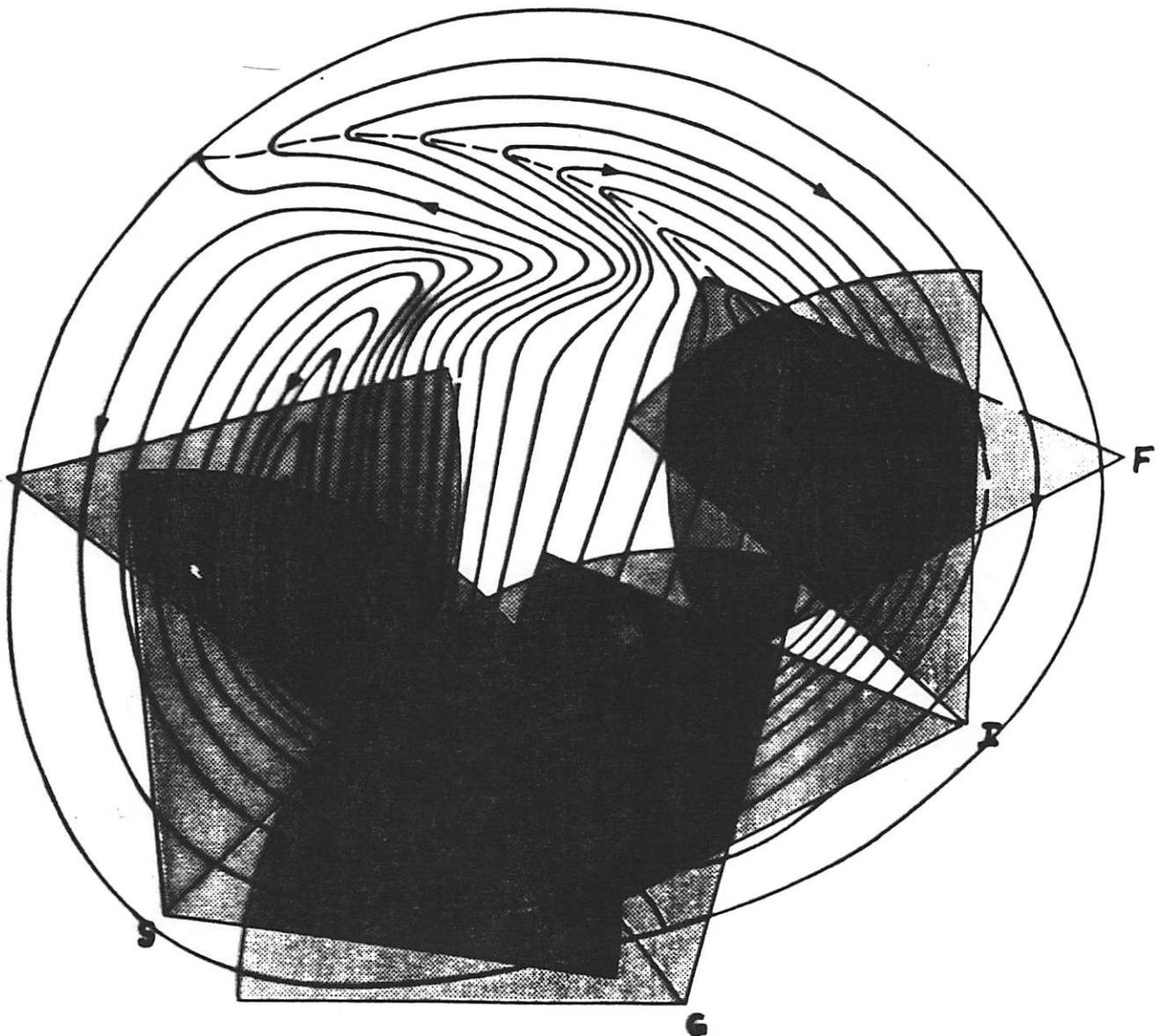
<u>Radar</u>	<u>Responsible Institutions</u>	<u>Status</u>
Finland Iceland (East)	University of Leicester (UK)	Funded
Iceland (West)	CNRS/LCPE, CRPE, LSEET (France) JHU/APL (US)	Funded
Goose Bay (E.)	JHU/APL	Operational
Goose Bay (W.)	JHU/APL (US) CNRS/LCPE, CRPE, LSEET (France)	Funded
Saskatoon (E)	U. of Saskatoon (Canada) also: U. of Alberta U. of Western Ont. U. of Calgary York University	Funded



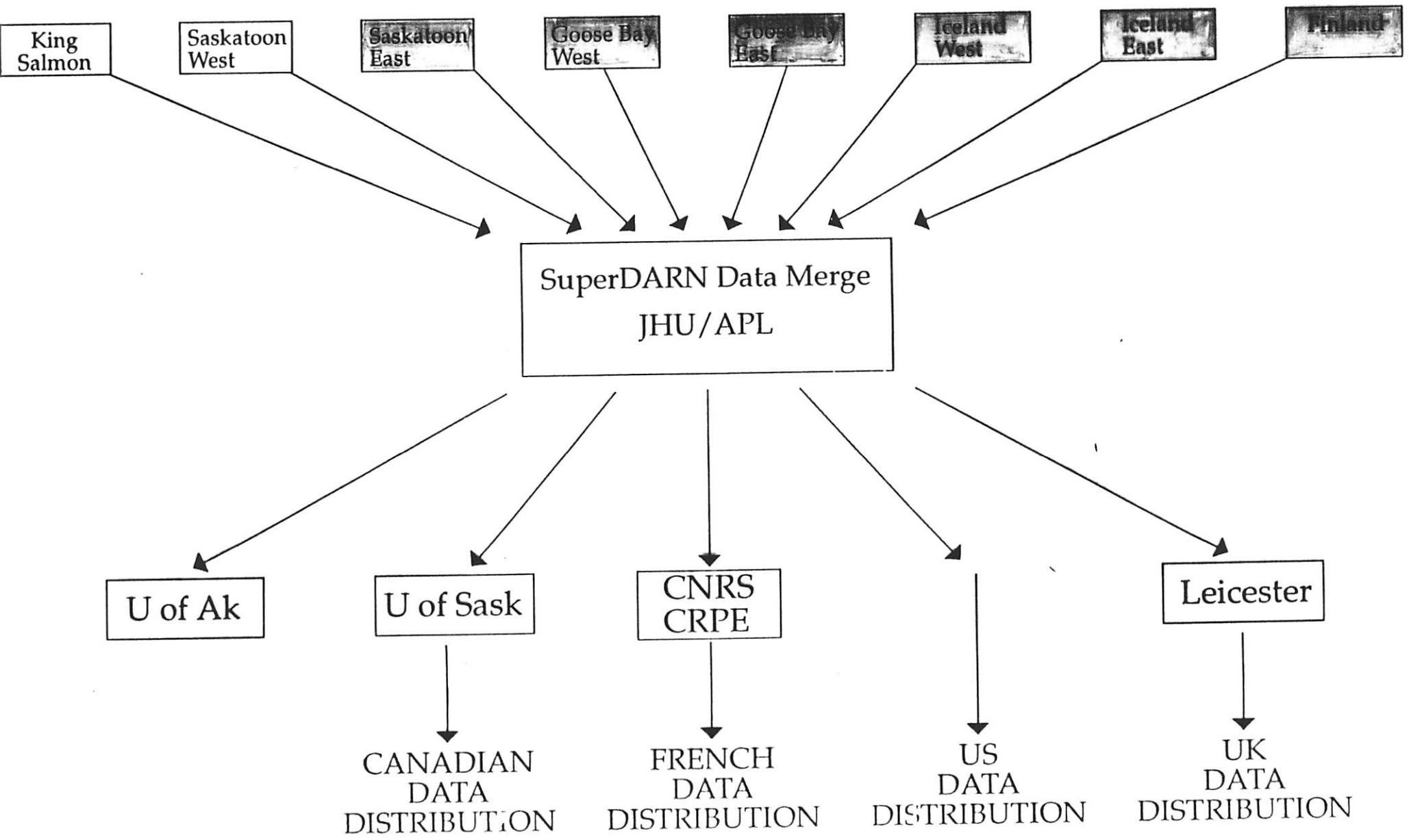
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SUPERDARN

Field-of-View Superimposed on Heppner-Maynard BC Pattern



SuperDARN Data Flow





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Anticipated Contributions of SuperDARN to Solar-Terrestrial Research

- **Data from SuperDARN will be shared by scientists in all participating countries**
- **Data from SuperDARN will be used to provide ground-based support to future multinational satellite missions including ISTP/GGS, CLUSTER, FREJA, and INTERBALL**
- **Data from SuperDARN will be used to support international scientific collaborations, e.g. STEP**
- **Data from SuperDARN will support national scientific initiatives, e.g. CEDAR, GEM**

